The Bramson correction in the Fisher-KPP equation: from delay to advance

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Abstract

We consider the solution to the scalar Fisher-KPP equation with front-like initial data, focusing on the location of its level sets at large times, particularly their deviation from points moving at the known spreading speed. We consider an intermediate case for the tail of the initial data, where the decay rate approaches, up to a polynomial term, that of the traveling wave with minimal speed. This approach enables us to capture deviations of the form $-r \ln t$ with $r < \frac{3}{2}$, which corresponds to a logarithmic delay when $0 < r < \frac{3}{2}$ and a logarithmic advance when r < 0. The critical case $r = \frac{3}{2}$ is also studied, revealing an extra $\mathcal{O}(\ln \ln t)$ term. Our arguments involve the construction of new sub- and super-solutions based on preliminary formal computations on the equation with a moving Dirichlet condition. Finally, convergence to the traveling wave with minimal speed is addressed.

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1 Introduction

This work is concerned with the large-time behavior of u = u(t, x), the solution to the Fisher-KPP equation

$$\partial_t u = \partial_{xx} u + u(1-u), \quad t > 0, \ x \in \mathbb{R}, \tag{1.1}$$

starting from a *front-like* initial data u_0 such that (see below for a precise statement)

$$u_0(x)$$
 "behaves like" $x^k e^{-x}$ as $x \to +\infty$, for some $k \ge -2$. (1.2)

As previously established, the level sets of the solution asymptotically move (as $t \to +\infty$) at a constant spreading speed, often denoted by c^* , which is also the minimal traveling wave speed. However, it is expected that the precise location of the level sets deviates slightly from the moving frame associated with this spreading speed. Through PDE techniques, we capture this (logarithmic) correction and, while the solution approaches a family of shifted minimal traveling waves, we find that it may either lag behind (when $-2 \le k < 1$) or slightly outpace (when k > 1) the moving frame with the exact speed c^* .

It is a well-established [4], [26] that, for large classes of front-like initial data (e.g. compactly supported or exponentially decaying as $x \to +\infty$), the solution of the Fisher-KPP equation (1.1) spreads with some speed $c \ge c^* = 2$, in the sense that

$$\min_{x \le c_1 t} u(t, x) \to 1 \text{ as } t \to +\infty \text{ if } c_1 < c, \quad \max_{x \ge c_2 t} u(t, x) \to 0 \text{ as } t \to +\infty \text{ if } c_2 > c.$$
(1.3)

The steepness of the right tail of the initial data significantly influences the position of the level sets. In particular, if the initial data has a "heavy tail", meaning it decays more slowly than any exponential as $x \to +\infty$, then the spreading speed is infinite and acceleration occurs [20].

Furthermore, in the case of exponential decay, and using a similar terminology to that of [28], one should distinguish between the *flat* case

$$u_0(x)$$
 "behaves like" $e^{-\lambda x}$ as $x \to +\infty$, for some $\lambda < 1$,

and the *sufficiently steep* case

$$u_0(x)$$
 "behaves like" $e^{-\lambda x}$ as $x \to +\infty$, for some $\lambda > 1$

In the flat case, the spreading speed is given by the dispersion relation $c = \frac{\lambda^2 + 1}{\lambda} > c^* = 2$, and there is no logarithmic correction. Specifically, the position of any level set of the solution behaves as $ct + \mathcal{O}(1)$ as $t \to +\infty$. We refer to [26], [27] for even stronger results on the convergence to a traveling wave. On the other hand, in the sufficiently steep case, the spreading speed is $c = c^* = 2$, but the position of any level set of the solution behaves as $2t - \frac{3}{2} \ln t + \mathcal{O}(1)$ as $t \to +\infty$, indicating a logarithmic delay. This was first known through probabilistic arguments [11, 12], and is commonly referred to as the logarithmic Bramson correction. As a result, the solution converges in the appropriate moving frame to the traveling

wave with minimal speed. This convergence can also be proved by other methods, e.g. based on an extended maximum principle related to the nonincrease of the number of intersection points between any two solutions of the Cauchy problem [22, 23]. Additionally, let us mention [13], [14], [28] for an analysis of such correction terms through convincing formal asymptotic expansions.

More recently, logarithmic corrections have been revisited through PDE techniques by Hamel, Nolen, Roquejoffre, and Ryzhik. In [19], they consider the case where $u_0 \equiv 0$ on a neighbourhood of $+\infty$ (roughly corresponding to $\lambda = +\infty$). However, a careful analysis of their proof reveals that it also applies to the case where

$$\left(x \mapsto x e^x u_0(x)\right) \in L^1(0, +\infty),$$

which, in particular, includes the sufficiently steep case. Further refinements, including the convergence rate, have then been developed by some of the same authors [24, 25]. In the same PDE framework, the connection between logarithmic corrections and the nature of waves (pushed vs. pulled) was explored in [16], [2, 3]. Additionally, the effect of nonlocal terms or discrete diffusion has recently drawn significant attention [9], [17], [10], [7]. Let us also mention the work [8] which shows that the delay can be increased by considering a logarithmic singular nonlinearity.

Our goal here is to investigate the intermediate case (1.2). The fact that the solution spreads with speed 2, in the sense that (1.3) holds with $c = c^* = 2$, follows directly from the previously mentioned results and the parabolic comparison principle, see also the insightful introduction in [18]. However, we will further prove, through PDE techniques akin to those in the aforementioned works, that such special initial conditions also lead to convergence to a traveling wave with a logarithmic correction, but with a prefactor depending (in value and sign) on k > -2. The critical case k = -2 will be tackled separately, as the position of the level sets in this case involves an extra term of order $\ln \ln t$.

Through this work, unless otherwise specified, we always denote by u = u(t, x) the solution of (1.1) with initial condition u_0 satisfying the following assumption.

Assumption 1.1 (Initial condition). The function $u_0 : \mathbb{R} \to [0,1]$ is uniformly continuous, positive, satisfies $\liminf_{x\to-\infty} u_0(x) > 0$ and there are $k \ge -2$, $0 < a \le A$ such that

$$ax^{k}e^{-x} \le u_{0}(x) \le Ax^{k}e^{-x}, \quad \forall x > 1.$$
 (1.4)

Notice that, by applying this assumption and the comparison principle, we immediately obtain 0 < u(t, x) < 1 for all $(t, x) \in (0, +\infty) \times \mathbb{R}$.

To state our results, we denote U the traveling wave with minimal speed (or simply minimal traveling wave) of (1.1), which satisfies the following conditions:

$$\begin{cases} U'' + 2U' + U(1 - U) = 0 \text{ on } \mathbb{R}, \\ U(-\infty) = 1, \ U(+\infty) = 0, \ U' < 0. \end{cases}$$
(1.5)

In particular U(x - 2t) solves (1.1). The existence and uniqueness (up to translation) of such a traveling wave were established in [15], [22], [4]. It is also known that

$$U(x) \sim Bxe^{-x}$$
, as $x \to +\infty$,

for some B > 0. In this context, when k = 1, respectively k < 1, respectively k > 1, Assumption 1.1 means that the right tail of the initial data roughly behaves like, is lighter than, or is heavier than that of the minimal traveling wave, respectively.

Our main result for the noncritical case k > -2 is as follows.

Theorem 1.2 (Bramson correction when k > -2). Let u = u(t, x) be the solution of (1.1) with initial data u_0 satisfying Assumption 1.1 with k > -2. Define

$$r := \frac{1-k}{2} < \frac{3}{2}.$$

Then, there exists a constant $C \ge 0$ such that

$$\lim_{t \to +\infty} \inf_{|h| \le C} \left\| u(t, \cdot) - U(\cdot - 2t + r \ln t + h) \right\|_{L^{\infty}(0, +\infty)} = 0.$$
(1.6)

Most of the proof will be devoted to analyzing the position of the level sets, from which the convergence part largely follows by a Liouville-type result as in [19]. More precisely, we will prove that, for any 0 < m < 1, there are $C_m, T_m > 0$ such that

$$E_m(t) \subset \left(2t - r\ln t - C_m, 2t - r\ln t + C_m\right), \text{ for all } t \ge T_m,$$

where

$$E_m(t) := \{ x \in \mathbb{R} : \ u(t, x) = m \}$$
(1.7)

denotes the *m*-level set of $u(t, \cdot)$ which, at least for large enough times, is not empty. As established in the aforementioned literature, the spreading result (1.3) holds with $c = c^* = 2$.

Let us briefly discuss the conclusions of Theorem 1.2.

- For the case k > 1, the Bramson correction implies that the precise location of the front of the solution moves faster than x = 2t. Therefore, we call it the *advance case*.
- Conversely, if -2 < k < 1, the Bramson correction implies that the precise location of the front of the solution moves slower than x = 2t. Hence, we refer to this as the *delay* case.
- When k = 1, there is no logarithmic deviation, which is consistent with the observation that, in this case, the right tail of the initial data is "comparable" to that of the traveling wave.

As for the critical case k = -2, we find that $r = \frac{1-k}{2} = \frac{3}{2}$. In this scenario, however, the drift is not expected to follow a purely logarithmic form. Instead, we provide the following result regarding the large-time asymptotics of the solution, which includes an extra $\mathcal{O}(\ln \ln t)$ term.

Theorem 1.3 (The critical case k = -2). Let u = u(t, x) be the solution to (1.1) starting from u_0 satisfying Assumption 1.1 with k = -2.

Then, there exists a constant $C \ge 0$ such that

t

$$\lim_{t \to +\infty} \inf_{|h| \le C} \left\| u(t, \cdot) - U\left(\cdot -2t + \frac{3}{2} \ln t - \ln \ln t + h \right) \right\|_{L^{\infty}(0, +\infty)} = 0.$$
(1.8)

We refer to the work [6], based on probabilistic arguments, for refined results in the critical case (but when the problem is instead posed on a moving half-line with a Dirichlet boundary condition). Let us also refer to the comprehensive work [14] which primarily utilizes formal asymptotic expansions. In contrast, our approach is distinct and relies on PDE techniques.

Outline of the paper. The organization of this work is as follows. In Section 2 we present some formal computations that are crucial to "guess" the shape of the solution, and provide material to construct sharp sub- and super-solutions. Building on these computations, we obtain the upper estimate on the level sets for the case k > -2 in Section 3 and the lower estimate in Section 4. The position of the level sets in the critical case k = -2 is studied in Section 5, using a similar but slightly different ansatz satisfying a moving Dirichlet boundary condition. Last, the convergence to the profile of the minimal traveling wave is addressed in Section 6.

2 Enlightening formal computations

The computations presented here are the keystone for constructing sharp sub- and supersolutions in the subsequent sections.

Computations involving the minimal wave. Denote as U the minimal Fisher-KPP wave, i.e. the (unique up to shifts) solution of (1.5). Then, defining $v(t, x) := U(x - 2t + r \ln(t + t_0))$ with $t_0 > 0$, we immediately compute

$$\mathcal{L}v := \partial_t v - \partial_{xx}v - v(1-v) = \frac{r}{t+t_0}U'\left(x - 2t + r\ln(t+t_0)\right).$$

When r > 0 (i.e. k = 1 - 2r < 1), then v acts as a sub-solution of (1.1). However, due to the asymptotics $U(x) \sim xe^{-x}$ as $x \to +\infty$, it follows that for any $t \ge 0$ that $v(t,x) > u_0(x)$ on some right half-line. Conversely, when r < 0 (i.e. k > 1), then v acts as a super-solution, but v(t,x) becomes smaller than $u_0(x)$ for large x. Consequently, in both cases, the comparison principle seems inapplicable.

Lemma 2.1. Let U be the minimal traveling wave and let $t_0 > 0$ be given. If $k \ge 1$, then $U(x - 2t + r \ln(t + t_0))$ is a super-solution of (1.1) on $[0, +\infty) \times \mathbb{R}$. If $k \le 1$, then $U(x - 2t + r \ln(t + t_0))$ is a sub-solution of (1.1) on $[0, +\infty) \times \mathbb{R}$.

This indicates that the traveling wave may serve in the so-called front zone $x \leq 2t + O(\sqrt{t})$. However, a different tool is required in the so-called far away zone $x \geq 2t + O(\sqrt{t})$. To do so, we work in the moving frame as in the seminal work [19], see also [16], leading us to an ODE problem. In contrast with previous literature, the adequate solution is determined by the choice of k > -2 in (1.4), which serves as a good starting point for constructing a relevant ansatz.

Computations in the moving frame. Following the idea from [19], we start from the linearized equation of (1.1) at $u \approx 0$,

$$\partial_t u = \partial_{xx} u + u.$$

By changing the variable

$$z = x - (2t - r\ln(t + t_0)),$$

we reach

$$\partial_t u = \partial_{zz} u + \left(2 - \frac{r}{t+t_0}\right) \partial_z u + u.$$

Letting

$$v(t,z) := e^z u(t,z),$$

we get

$$\partial_t v = \partial_{zz} v + \frac{r}{t+t_0} (v - \partial_z v).$$

Using the self-similar variables

$$\tau = \ln(t + t_0) - \ln t_0, \quad y = \frac{z}{\sqrt{t + t_0}},$$

we reach

$$\begin{aligned} \partial_{\tau} v &= \partial_{yy} v + \frac{y}{2} \partial_{y} v + rv - \frac{r}{\sqrt{t_0}} e^{-\frac{\tau}{2}} \partial_{y} v \\ &\approx \partial_{yy} v + \frac{y}{2} \partial_{y} v + rv, \end{aligned}$$

by formally ignoring the non-autonomous exponentially decreasing drift term. Plugging the ansatz

$$v(\tau, y) = e^{\frac{\tau}{2}}w(y),$$

with w(0) = 0 and w'(0) = 1 (which can be motivated by a matching argument with the traveling wave), we are left to solve the Cauchy problem

$$\begin{cases} w'' + \frac{y}{2}w' + \left(r - \frac{1}{2}\right)w = 0, \\ w(0) = 0, \quad w'(0) = 1. \end{cases}$$

Lemma 2.2. The solution w of the Cauchy problem

$$\begin{cases} w'' + \frac{y}{2}w' + \left(r - \frac{1}{2}\right)w = 0, \quad y > 0, \\ w(0) = 0, \\ w'(0) = 1, \end{cases}$$
(2.1)

satisfies the following properties:

(i) if
$$r = \frac{3}{2}$$
, then $w(y) = ye^{-y^2/4}$;
(ii) if $r < \frac{3}{2}$, then $w > 0$ in $(0, +\infty)$, $\lim_{y \to +\infty} \frac{w'}{w}(y) = 0$ and there exists $C > 0$ such that
 $w(y) \sim Cy^k$ as $y \to +\infty$, $k = 1 - 2r$. (2.2)

Proof. The case $r = \frac{3}{2}$ is obvious as one may directly check that $ye^{-y^2/4}$ indeed satisfies the wanted problem. When $r < \frac{3}{2}$, we may write $w(y) = y\varphi(y)$ where

$$\begin{cases} \varphi'' + \left(\frac{y}{2} + \frac{2}{y}\right)\varphi' + r\varphi = 0, \quad y > 0, \\ \varphi(0) = 1, \\ \varphi'(0) = 0. \end{cases}$$

Notice that, using a Sturm-Liouville approach, one can recast the above ODE problem into an integral equation and prove the existence and uniqueness of a local solution which satisfies $\varphi''(0) = -\frac{r}{3}$, see [21, Proposition 3.1]. Next, by writing $\varphi(y) = e^{-\frac{y^2}{4}}\psi(\frac{y^2}{4})$, we see that ψ satisfies

$$\begin{cases} z\psi'' + \left(\frac{3}{2} - z\right)\psi' + \left(r - \frac{3}{2}\right)\psi = 0, \quad z > 0\\ \psi(0) = 1, \\ \psi'(0) = 1 - \frac{2r}{3}. \end{cases}$$

Then $\psi(z) = {}_{1}F_{1}\left(\frac{3-2r}{2}, \frac{3}{2}, z\right)$, where ${}_{1}F_{1}(a, b, z)$ denotes the confluent hypergeometric function of the first kind, or Kummer's function [1]. According to [1, formula 13.1.4], for a not a nonpositive integer, we have

$$_{1}F_{1}(a,b,z) \sim \frac{\Gamma(b)}{\Gamma(a)} \frac{e^{z}}{z^{b-a}}, \quad \text{as } z \to +\infty,$$

which, applied to $w(y) = ye^{-\frac{y^2}{4}}\psi\left(\frac{y^2}{4}\right)$, exactly corresponds to (2.2). Additionally, from [1, formula 13.4.14], we have

$$\sqrt{z} \left(\frac{\frac{d}{dz_1} F_1(a, b, z)}{{}_1F_1(a, b, z)} - 1 \right) \to 0, \quad \text{as } z \to +\infty,$$

which implies $\frac{w'(y)}{w(y)} \to 0$ as $y \to +\infty$. Lastly, since $a = \frac{3-2r}{2} > 0$, $b = \frac{3}{2}$, it follows from [1, formula 13.1.2] that ψ is positive on $(0, +\infty)$ and so is w.

Observe that, returning to u = u(t, z), we have

$$u(t,z) = e^{-z} \sqrt{\frac{t+t_0}{t_0}} w\left(\frac{z}{\sqrt{t+t_0}}\right),$$

so that not only the asymptotic behavior (2.2) ensures that

$$u(0,z) \sim \frac{C}{t_0^{k/2}} z^k e^{-z}, \quad \text{as } z \to +\infty,$$

which matches with the spatial decay of the initial data as stated in Assumption 1.1, but also ensures

$$u(t,z) \sim \frac{1}{\sqrt{t_0}} z e^{-z}, \quad \text{as } t \to +\infty,$$

which matches with the asymptotics of the minimal traveling wave. This suggests that this ansatz may serve effectively as either a sub-solution or a super-solution.

Remark 2.3. Incorporating the nonlinear and non-autonomous terms that were previously ignored in the formal computation, the proposed ansatz, unfortunately, fails to generate either a valid sub-solution or super-solution. To overcome this issue, we will consider a slightly different approach, introducing a logarithmic drift different from the expected one. More precisely, taking r' > 0 and letting instead $z = x - (2t - r' \ln(t + t_0))$ and

$$v(\tau, y) = e^{(r'-r)\tau} e^{\frac{\tau}{2}} w(y)$$

in the above computation, we observe that w still satisfies the ODE (2.1). Thanks to this additional parameter r', the resulting alternative ansatz now satisfies a Dirichlet boundary condition at $x = 2t + r' \ln(t + t_0)$, and still exhibits similar asymptotic behavior. However, the inclusion of r' proves to be more convenient to construct sub- and super-solutions of (1.1).

3 The upper estimate

This section focuses on establishing the upper estimate for the position of the *m*-level set E_m of the solution, as defined in (1.7).

Proposition 3.1 (Upper estimate). Let u = u(t, x) be the solution to (1.1) starting from u_0 satisfying Assumption 1.1 with k > -2. Let $m \in (0, 1)$ be given. Then, there are $C_m > 0$ and $T_m > 0$ such that

$$E_m(t) \subset (-\infty, 2t - r \ln t + C_m), \quad \forall t \ge T_m,$$
(3.1)

with

$$r = \frac{1-k}{2} < \frac{3}{2}.$$

Moreover, there are $k_1 > 0$ and $\sigma_1 > 0$ such that

$$u(t, 2t - r \ln t + y) \le k_1(y+1)e^{-y}$$
 for any $t \ge 1, \ 0 \le y \le \sigma_1 \sqrt{t}$. (3.2)

3.1 Proof of Proposition 3.1

To construct an appropriate ansatz for the solution, we begin by considering the form derived from the formal computations in the moving frame (see Section 2), which suggests the following behavior:

$$e^{-z}\sqrt{t}w\left(\frac{z}{\sqrt{t}}\right), \quad z = x - 2t + r\ln t,$$

with w coming from Lemma 2.2. However, some modifications will be necessary to deal, in particular, with the non-autonomous drift term that was neglected in Section 2, see Remark 2.3. However, this ansatz ignores the non-autonomous drift term that we neglected earlier in the derivation, and incorporating this term will be crucial to constructing proper sub- and super-solutions. These modifications slightly differ between the delay case and the advance case, as the signs of those ignored terms are different.

Step 1: a first super-solution. For $t > 0, z \ge 0$, we define

$$\psi(t,z) := e^{-z} t^{\frac{1}{2} + r' - r} w\left(\frac{z}{\sqrt{t}}\right), \qquad (3.3)$$

where the function w is given by Lemma 2.2, and

$$r' = \max(r, 0). \tag{3.4}$$

For $M \ge 0$ to be specified later, let us prove that

$$\widehat{u}(t,x) := \left(1 - \frac{M}{\sqrt{t}}\right)\psi(t,x - 2t + r'\ln t)$$
(3.5)

is a super-solution of (1.1) for all $t \ge t_0$ (with $t_0 > 1$ large enough) and $x \ge 2t - r' \ln t$.

Let us start with

$$\tilde{u}(t,x) := \psi(t, x - 2t + r' \ln t).$$

Observe that

$$\begin{aligned} \partial_t \psi(t,z) &= \left(\frac{1}{2} + r' - r\right) \frac{e^{-z}}{t^{\frac{1}{2} + r - r'}} w\left(\frac{z}{\sqrt{t}}\right) - \frac{ze^{-z}}{2t^{1 + r - r'}} w'\left(\frac{z}{\sqrt{t}}\right), \\ \partial_z \psi(t,z) &= -\psi(t,z) + e^{-z} t^{r' - r} w'\left(\frac{z}{\sqrt{t}}\right), \\ \partial_{zz} \psi(t,z) &= \psi(t,z) - 2e^{-z} t^{r' - r} w'\left(\frac{z}{\sqrt{t}}\right) + \frac{e^{-z}}{t^{\frac{1}{2} + r - r'}} w''\left(\frac{z}{\sqrt{t}}\right). \end{aligned}$$

Combining this and (2.1), we reach

$$\begin{split} \mathcal{L}\tilde{u}(t,x) &\geq \partial_{t}\tilde{u} - \partial_{xx}\tilde{u} - \tilde{u} \\ &= \partial_{t}\psi + \left(-2 + \frac{r'}{t}\right)\partial_{z}\psi - \partial_{zz}\psi - \psi \\ &= e^{-z}\left(\frac{\frac{1}{2} + r' - r}{t^{\frac{1}{2} + r - r'}} - \frac{r'}{t^{\frac{1}{2} + r - r'}}\right)w\left(\frac{z}{\sqrt{t}}\right) + e^{-z}\left(-\frac{z}{2t^{1 + r - r'}} + \frac{r'}{t^{1 + r - r'}}\right)w'\left(\frac{z}{\sqrt{t}}\right) \\ &\quad -\frac{e^{-z}}{t^{\frac{1}{2} + r - r'}}w''\left(\frac{z}{\sqrt{t}}\right) \\ &= e^{-z}\frac{\frac{1}{2} - r}{t^{\frac{1}{2} + r - r'}}w\left(\frac{z}{\sqrt{t}}\right) + e^{-z}\left(-\frac{z}{2t^{1 + r - r'}} + \frac{r'}{t^{1 + r - r'}}\right)w'\left(\frac{z}{\sqrt{t}}\right) \\ &\quad +\frac{e^{-z}}{t^{\frac{1}{2} + r - r'}}\left(\frac{z}{2\sqrt{t}}w'\left(\frac{z}{\sqrt{t}}\right) + \left(r - \frac{1}{2}\right)w\left(\frac{z}{\sqrt{t}}\right)\right) \\ &= \frac{r'e^{-z}}{t^{1 + r - r'}}w'\left(\frac{z}{\sqrt{t}}\right), \end{split}$$

where we use the shorthand $z = x - 2t + r' \ln t$.

When $r \leq 0$, then r' = 0 and \tilde{u} is already a super-solution, so one may choose M = 0in (3.5). Unfortunately, when r > 0 this remaining term may be negative, and at least this is the case when $r > \frac{1}{2}$ due to $w(y) \sim Cy^{1-2r}$ as $y \to +\infty$. We now deal with the case r > 0 (so that r' = r). The key point is that w is increasing

We now deal with the case r > 0 (so that r' = r). The key point is that w is increasing on a neighborhood of 0, beyond which |w'|/w remains bounded. Precisely, from Lemma 2.2, we can select $\delta > 0$ such that w' > 0 in $[0, \delta)$, and then M > 0 such that

$$\frac{2r'|w'|}{w} \le M \quad \text{on} \quad [\delta, +\infty). \tag{3.6}$$

Now, we set $t_0 > 1$ large enough such that $1 - \frac{M}{\sqrt{t_0}} \ge \frac{1}{2}$. Then, since $\hat{u}(t, x) = \left(1 - \frac{M}{\sqrt{t}}\right) \tilde{u}(t, x)$, one has

$$\mathcal{L}\widehat{u} \ge \partial_t \widehat{u} - \partial_{xx}\widehat{u} - \widehat{u} = \frac{e^{-z}}{t} \left[\left(1 - \frac{M}{\sqrt{t}} \right) r'w' \left(\frac{z}{\sqrt{t}} \right) + \frac{M}{2} w \left(\frac{z}{\sqrt{t}} \right) \right] \ge 0,$$

for any $t \ge t_0$, $z = x - 2t + r' \ln t \ge 0$. The last inequality follows from either $w'\left(\frac{z}{\sqrt{t}}\right) > 0$ for $0 \le z \le \delta\sqrt{t}$, or from (3.6) for $z \ge \delta\sqrt{t}$. We have thus proved that $\hat{u} = \hat{u}(t,x)$ is a super-solution of (1.1) for all $t \ge t_0$ and $x \ge 2t - r' \ln t$. Step 2: comparison with the solution. Since \hat{u} from (3.5) satisfies $\hat{u}(t, 2t - r' \ln t) = 0$ for $t \ge 0$, we cannot directly use it as a super-solution. Therefore we will first cut it by 1 on a left half-line. By setting $t_0 > 1$ large enough such that $\frac{M}{\sqrt{t_0}} \le \frac{1}{2}$, we have

$$\begin{aligned} \widehat{u} \left(t + t_0, 2(t + t_0) - r \ln(t + t_0) + 1 \right) \\ &= \left(1 - \frac{M}{\sqrt{t + t_0}} \right) e^{-1 - (r' - r) \ln(t + t_0)} \left(t + t_0 \right)^{\frac{1}{2} + r' - r} w \left(\frac{(r' - r) \ln(t + t_0) + 1}{\sqrt{t + t_0}} \right) \\ &\geq \frac{e^{-1}}{2} \sqrt{t + t_0} \times w \left(\frac{(r' - r) \ln(t + t_0) + 1}{\sqrt{t + t_0}} \right). \end{aligned}$$

We deduce from w'(0) = 1 that

$$w\left(\frac{(r'-r)\ln(t+t_0)+1}{\sqrt{t+t_0}}\right) \sim \frac{(r'-r)\ln(t+t_0)+1}{\sqrt{t+t_0}} \quad \text{as } t \to +\infty$$

hence (using also the fact that $r' - r \ge 0$)

$$\inf_{t \ge 0} \ \widehat{u} \left(t + t_0, 2(t + t_0) - r \ln(t + t_0) + 1 \right) > 0$$

It follows that we can find $K\gg 1$ so that

$$K\widehat{u}\left(t+t_{0}, 2(t+t_{0})-r\ln(t+t_{0})+1\right) > 1,$$

for all $t \ge 0$. Then

$$\overline{u}(t,x) := \begin{cases} 1 & \text{if } x < 2(t+t_0) - r\ln(t+t_0) + 1, \\ \min\{1, K\widehat{u}(t+t_0, x)\} & \text{if } x \ge 2(t+t_0) - r\ln(t+t_0) + 1, \end{cases}$$

is a generalized super-solution of (1.1).

Next we prove that, up to increasing K > 0 if necessary,

$$u_0(x) \le \overline{u}(0, x). \tag{3.7}$$

Since $u_0 \leq 1$, we only have to check that

$$\forall x \ge 2t_0 - r \ln t_0 + 1, \quad u_0(x) \le K \widehat{u}(t_0, x). \tag{3.8}$$

By (3.3), (3.5), we have

$$\widehat{u}(t_0, x) \ge C(t_0)e^{-x}w\left(\frac{x - 2t_0 + r'\ln t_0}{\sqrt{t_0}}\right)$$

for some $C(t_0) > 0$. Then, due to $w(y) \sim Cy^k$ as $y \to +\infty$ with C > 0 according to Lemma 2.2, up to decreasing the value of the constant $C(t_0)$ we get that

$$\forall x \ge 2t_0 - r \ln t_0 + 1, \quad \widehat{u}(t_0, x) \ge C(t_0) x^k e^{-x}.$$

From this and (1.4) it follows that, for K > 0 large enough, (3.8) holds, and so does (3.7). However, we point out that the assumption $r < \frac{3}{2}$ played a crucial role here, since when $r = \frac{3}{2}$ then w(y) has a gaussian behavior as $y \to +\infty$ and (3.8) cannot hold regardless of the multiplicative constant K.

Still, when $r < \frac{3}{2}$, and applying the parabolic comparison principle, we get that

$$u(t,x) \le \overline{u}(t,x),\tag{3.9}$$

for any $t \ge 0$ and $x \in \mathbb{R}$.

In particular, for t large enough and $x \ge 2(t+t_0) - r \ln(t+t_0) + \sqrt{t}$, we find that

$$\begin{split} u(t,x) &\leq K e^{-(x-2(t+t_0)+r'\ln(t+t_0))}(t+t_0)^{\frac{1}{2}+r'-r} w \left(\frac{x-2(t+t_0)+r'\ln(t+t_0)}{\sqrt{t+t_0}}\right) \\ &\leq K \left(\sup_{y\geq 1} \frac{w(y)}{y^k}\right) e^{-(x-2(t+t_0))}(t+t_0)^{\frac{1}{2}-r} \times \left(\frac{x-2(t+t_0)+r'\ln(t+t_0)}{\sqrt{t+t_0}}\right)^k \\ &\leq K \left(\sup_{y\geq 1} \frac{w(y)}{y^k}\right) e^{-(x-2(t+t_0))} \left(x-2(t+t_0)+r'\ln(t+t_0)\right)^k \\ &\leq K \left(\sup_{y\geq 1} \frac{w(y)}{y^k}\right) e^{-(\sqrt{t}-r\ln(t+t_0))}(\sqrt{t}+(r'-r)\ln(t+t_0))^k, \end{split}$$

where we have used the fact that k = 1 - 2r. Since Lemma 2.2 ensures that $\sup_{y \ge 1} \frac{w(y)}{y^k} < +\infty$, we conclude that

$$\lim_{t \to +\infty} \sup_{x \ge 2(t+t_0) - r \ln(t+t_0) + \sqrt{t}} u(t, x) = 0.$$
(3.10)

To obtain an upper estimate between the expected moving frame of the level sets, i.e. $2t - r \ln t + \mathcal{O}(1)$, and $2(t+t_0) - r \ln(t+t_0) + \sqrt{t}$ beyond which the above limit holds, the proof differs slightly between the delay and advance cases. We detail this in the next two steps.

Step 3: conclusion in the delay case $0 \le r < \frac{3}{2}$. Here, r' = r by (3.4) so that (3.9) yields

$$u(t,x) \le K e^{-(x-2(t+t_0)+r\ln(t+t_0))} \sqrt{t+t_0} \times w\left(\frac{x-2(t+t_0)+r\ln(t+t_0)}{\sqrt{t+t_0}}\right),$$

for $x \ge 2(t+t_0)-r\ln(t+t_0)+1$, where $t_0 > 1$ is fixed large enough. Recalling that $w(y) \sim Cy^k$ with $k = 1 - 2r \le 1$ as $y \to +\infty$, and also that w(0) = 0 and w'(0) = 1, we have that w(y)/y is bounded on $(0, +\infty)$. Then

$$u(t,x) \le K\left(\sup_{y\ge 0}\frac{w(y)}{y}\right)(x-2(t+t_0)+r\ln(t+t_0))e^{-(x-2(t+t_0)+r\ln(t+t_0))}$$

for $x \ge 2(t+t_0) - r \ln(t+t_0) + 1$, from which the conclusion (3.1) of Proposition 3.1 follows. Furthermore, evaluating the previous inequality at $x = 2t - r \ln t + y$ with $t \ge 1$ and $y \ge 2t_0 + 1$, we find some $k_1 > 0$ large enough so that

$$u(t, 2t - r \ln t + y) \le k_1(y+1)e^{-y}$$
 for any $t \ge 1, y \ge 2t_0 + 1$.

From the global boundedness of u and the positivity of the right-hand term for $y \ge 0$, one then infers (3.2), which concludes the proof of Proposition 3.1 in the delay case.

Step 4: conclusion in the advance case r < 0. Here, r' = 0 > r. By (3.4), the same computation as in Step 3 only provides an upper estimate on the location of the level sets of the form $2t - r \ln t + C_m \ln(\ln t)$. Therefore we must refine the argument. First, for $t \ge 1$, (3.9) now yields

$$u\left(t, 2(t+t_0) - r\ln(t+t_0) + \sqrt{t}\right) \le Ke^{-\sqrt{t}}\sqrt{t+t_0} \times w\left(\frac{-r\ln(t+t_0) + \sqrt{t}}{\sqrt{t+t_0}}\right),$$

and thus there is $T \gg 1$ such that

$$u\left(t, 2(t+t_0) - r\ln(t+t_0) + \sqrt{t}\right) \le 2K\sqrt{t}e^{-\sqrt{t}}w(1), \quad \forall t \ge T.$$
(3.11)

Besides, since

$$\varphi(z) := z e^{-z}$$

satisfies $\varphi'' + 2\varphi' + \varphi = 0$ and $\varphi' \leq 0$ on $(1, +\infty)$, and recalling that r < 0 here, it is straightforward that, for any $K_2 > e$,

$$\overline{u}_2(t,x) := \begin{cases} 1 & \text{if } x < 2t - r \ln t + 1, \\ \min\{1, K_2 \,\varphi(x - 2t + r \ln t)\} & \text{if } x \ge 2t - r \ln t + 1, \end{cases}$$

is a generalized super-solution of (1.1). Furthermore,

$$\overline{u}_2(t, 2(t+t_0) - r\ln(t+t_0) + \sqrt{t}) \sim K_2 e^{-2t_0} \sqrt{t} e^{-\sqrt{t}}$$
 as $t \to +\infty$.

Putting this together with (3.11), we get, up to increasing T if necessary, that

$$u\left(t, 2(t+t_0) - r\ln(t+t_0) + \sqrt{t}\right) \le \overline{u}_2\left(t, 2(t+t_0) - r\ln(t+t_0) + \sqrt{t}\right), \quad \forall t \ge T,$$

provided that $K_2 > 2Kw(1)e^{2t_0}$. As in Step 2, due to $u \leq 1$, we may increase $K_2 > 0$ if necessary so that

$$u(T,x) \le \overline{u}_2(T,x)$$

for all $x \leq 2(T+t_0) - r \ln(T+t_0) + \sqrt{T}$. Applying the parabolic comparison principle on

$$\Omega_T := \{(t, x) : t \ge T \text{ and } x \le 2(t + t_0) - r \ln(t + t_0) + \sqrt{t}\},\$$

we end up with

$$u(t,x) \leq \overline{u}_2(t,x), \quad \forall (t,x) \in \Omega_T.$$

The estimate (3.2) follows immediately from the definition of \overline{u}_2 , at least for $t \geq T^*$ with $T^* > 1$ large enough and $1 \leq y \leq \frac{\sqrt{t}}{2}$ with $k_1 = K_2$. For $0 \leq y \leq 1$, or for $1 \leq t \leq T^*$ and $0 \leq y \leq \frac{\sqrt{t}}{2}$, the same estimate holds up to increasing the factor k_1 , thanks to the boundedness of u. Moreover, putting the previous inequality together with (3.10), since any m-level set of \overline{u}_2 is located around $2t - r \ln t + \mathcal{O}(1)$, we finally deduce (3.1).

3.2 An alternate proof in the advance case (k > 1)

We start with an estimate on the solution v(t, x) to (3.12) (the linearized equation of (1.1) around the invading steady state) starting from u_0 , at a moving point $x = 2t + \mathcal{O}(\sqrt{t})$.

Lemma 3.2. Let v = v(t, x) be the solution to

$$\partial_t v = \partial_{xx} v + v, \quad t > 0, \ x \in \mathbb{R}, \tag{3.12}$$

starting from u_0 satisfying Assumption 1.1 with k > 1. Let c > 0 and $t_0 > 0$ be given. Then there is $c_0 > 0$ such that

$$v(t, 2t + c\sqrt{t}) \le c_0 t^{\frac{k}{2}} e^{-c\sqrt{t}}, \quad \forall t \ge t_0.$$

Proof. Denoting $G(t,x) := \frac{1}{\sqrt{4\pi t}}e^{-\frac{x^2}{4t}}$ the one dimensional Heat kernel, we have

$$e^{-t}v(t,x) = \int_{\mathbb{R}} G(t,x-y)u_0(y)\,dy \le \int_{-\infty}^1 G(t,x-y)\,dy + A\int_1^{+\infty} G(t,x-y)y^k e^{-y}\,dy, \quad (3.13)$$

from Assumption 1.1. Evaluating at $x = 2t + c\sqrt{t}$, the first term in the above right hand side is nothing else than

$$\frac{1}{\sqrt{\pi}} \int_{-\infty}^{-\sqrt{t} - \frac{c}{2} + \frac{1}{2\sqrt{t}}} e^{-z^2} dz \sim \frac{1}{2\sqrt{\pi}} \frac{1}{\sqrt{t}} e^{-(\sqrt{t} + \frac{c}{2} - \frac{1}{2\sqrt{t}})^2}, \text{ as } t \to +\infty.$$
(3.14)

As for the second term, using the change of variable $y = 2\sqrt{t}z + x - 2t$ with $x = 2t + c\sqrt{t}$, some straightforward computations yield

$$A\int_{1}^{+\infty} G(t, x-y)y^{k}e^{-y}\,dy = A\frac{t^{\frac{k}{2}}e^{-(t+c\sqrt{t})}}{\sqrt{\pi}}\int_{-\frac{c}{2}+\frac{1}{2\sqrt{t}}}^{+\infty} (2z+c)^{k}e^{-z^{2}}\,dz \le c't^{\frac{k}{2}}e^{-(t+c\sqrt{t})},$$
(3.15)

for some c' > 0. The combination of (3.13), (3.14) and (3.15) concludes the proof (and so if we only assume k > -1 rather than k > 1).

The above upper estimate now allows us to use the minimal traveling wave, see Lemma 2.1, as an accurate super-solution.

Proof of Proposition 3.1 when k > 1. We start with some preparation. From Lemma 3.2 (and the comparison principle), there is $c_0 > 0$ such that

$$u(t, 2t + \sqrt{t}) \le c_0 t^{\frac{k}{2}} e^{-\sqrt{t}}, \quad \forall t \ge 1.$$

Since $U(z) \sim ze^{-z}$ as $z \to +\infty$, we can find $z_0 > 0$ such that $U(z) \ge \frac{1}{2}ze^{-z}$ for all $z \ge z_0$. Select T > 1 large enough such that

$$\sqrt{t} + r \ln(1+t) \ge \frac{1}{2}\sqrt{t} \ge z_0, \quad \forall t \ge T,$$

and K > 1 large enough such that

$$K > 4c_0$$
, and $KU(\sqrt{T} + r\ln(1+T)) > 1.$ (3.16)

We aim at applying the comparison principle in the region $\Omega_T := \{(t, x) \in [T, +\infty) \times \mathbb{R} : x \leq 2t + \sqrt{t}\}$. We define

$$\overline{u}(t,x) := KU(x - 2t + r\ln(1+t)),$$

which is a super-solution since

$$\partial_t \overline{u} - \partial_{xx} \overline{u} - \overline{u}(1 - \overline{u}) = K \left(-2 + \frac{r}{1+t} \right) U' - K U'' - K U + K^2 U^2$$
$$= K \frac{r}{1+T} U' + (K^2 - K) U^2$$

is positive since K > 1, r < 0, U' < 0. On the lateral boundary $(t \ge T, x = 2t + \sqrt{t})$ of Ω_T , we have

$$\begin{split} \overline{u}(t, 2t + \sqrt{t}) &= KU(\sqrt{t} + r\ln(1+t)) \\ &\geq \frac{K}{2}(\sqrt{t} + r\ln(1+t))e^{-\sqrt{t} - r\ln(1+t)} \\ &\geq \frac{K}{4}\sqrt{t}e^{-\sqrt{t}}(1+t)^{-r} \geq \frac{K}{4}t^{\frac{k}{2}}e^{-\sqrt{t}} \geq u(t, 2t + \sqrt{t}), \end{split}$$

from (3.16). On the boundary $(t = T, x \leq 2T + \sqrt{T})$ of Ω_T we have

$$\overline{u}(T,x) \ge KU(\sqrt{T} + r\ln(1+T)) > 1 \ge u(T,x).$$

We thus deduce from the comparison principle that

$$u(t,x) \leq KU(x-2t+r\ln(1+t)), \quad \text{for all } (t,x) \in \Omega_T,$$

from which the conclusion of Proposition 3.1 easily follows.

4 The lower estimate

This section is devoted to the following lower estimate. Recall that E_m denotes the *m*-level set of the solution, as defined in (1.7).

Proposition 4.1 (Lower estimate). Let u = u(t, x) be the solution to (1.1) starting from u_0 satisfying Assumption 1.1 with k > -2. Let $m \in (0, 1)$ be given. Then, there are $C_m > 0$ and $T_m > 0$ such that

$$E_m(t) \subset (2t - r \ln t - C_m, +\infty), \quad \forall t \ge T_m,$$

$$(4.1)$$

where

$$r = \frac{1-k}{2} < \frac{3}{2}.$$

Moreover, there are $k_2 > 0$ and $\sigma_2 > 0$ such that

$$u(t, 2t - r \ln t + y) \ge k_2 y e^{-y}$$
 for any $t \ge 1, \ 0 \le y \le \sigma_2 \sqrt{t}$. (4.2)

4.1 A lower bound in the far away zone

We begin by demonstrating a lower estimate on the solution to (1.1), starting from u_0 , evaluated at a moving point $x = 2t + \mathcal{O}(\sqrt{t})$. Note that this estimate aligns with a shift of the minimal front U, so that its level sets are positioned around $x = 2t - r \ln t$. This will turn out to be crucial for bootstrapping our argument with a second sub-solution in subsection 4.2.

Lemma 4.2. Let u = u(t, x) be the solution to (1.1) starting from u_0 satisfying Assumption 1.1 with k > -2. Then there exists $\varepsilon > 0$ such that

$$u(t, 2t + \sqrt{t}) \ge \varepsilon e^{-\sqrt{t}} t^{\frac{1}{2}-r}, \quad \forall t \ge T_0,$$

$$(4.3)$$

for some $T_0 > 1$ large enough.

Proof. We proceed into two steps. The first one consists in turning our computations from Section 2 into an actual sub-solution of (1.1). Next, we will prove that this sub-solution provides the lower estimate at $x = 2t + \sqrt{t}$.

Step 1: the first sub-solution. For $t > 0, z \ge 0$, we define as before

$$\psi(t,z) := e^{-z} t^{\frac{1}{2} + r' - r} w\left(\frac{z}{\sqrt{t}}\right), \tag{4.4}$$

where the function w is given by Lemma 2.2, and now

$$r' = r - 2 < 0. \tag{4.5}$$

For $t_0 > 1$, M > 0, $\varepsilon > 0$ to be specified later, let us prove that

$$\underline{u}(t,x) := \left(1 + \frac{M}{\sqrt{t}}\right) \varepsilon \,\psi(t,x - 2t + r' \ln t) \tag{4.6}$$

is a sub-solution of (1.1) for all $t \ge t_0$ and $x \ge 2t - r' \ln t$. By construction $\underline{u}(t, x)$ also satisfies the Dirichlet boundary condition

$$\underline{u}(t, 2t - r'\ln t) = 0.$$

We use the computations of Step 1 in subsection 3.1 again and get, denoting $z = x - 2t + r' \ln t$,

$$\begin{aligned} \mathcal{L}\underline{u}(t,x) &= \partial_t \underline{u} - \partial_{xx} \underline{u} - \underline{u} + \underline{u}^2 \\ &= \varepsilon \frac{e^{-z}}{t^{1+r-r'}} \Biggl\{ -\frac{M}{2} w \left(\frac{z}{\sqrt{t}} \right) \\ &+ \left(1 + \frac{M}{\sqrt{t}} \right) \left[r'w' \left(\frac{z}{\sqrt{t}} \right) + \left(1 + \frac{M}{\sqrt{t}} \right) \varepsilon e^{-z} t^{2+r'-r} w^2 \left(\frac{z}{\sqrt{t}} \right) \right] \Biggr\} \\ &= \varepsilon \frac{e^{-z}}{t^3} \Biggl\{ -\frac{M}{2} w \left(\frac{z}{\sqrt{t}} \right) + \left(1 + \frac{M}{\sqrt{t}} \right) \left[r'w' \left(\frac{z}{\sqrt{t}} \right) + \left(1 + \frac{M}{\sqrt{t}} \right) \varepsilon e^{-z} w^2 \left(\frac{z}{\sqrt{t}} \right) \right] \Biggr\}, \end{aligned}$$

where on the last line we simply used r' = r - 2.

Next, we distinguish two subdomains. First, we use the fact that w'(0) = 1, hence there exists $\delta > 0$ such that $w'(y) > \frac{1}{2}$ in $y \in [0, \delta]$. Then, due to r' < 0, we have for any $t \ge 1$ and $z = x - 2t + r' \ln t \in (0, \delta\sqrt{t}]$,

$$\mathcal{L}\underline{u}(t,x) \leq \varepsilon \frac{e^{-z}}{t^3} \left(1 + \frac{M}{\sqrt{t}} \right) \left[\frac{r'}{2} + (1+M) \varepsilon \max_{0 \leq y \leq \delta} w^2(y) \right],$$

which is negative for any

$$\varepsilon < \frac{-r'}{2(1+M)\max_{0 \le y \le \delta} w^2(y)}.$$
(4.7)

Furthermore, from Lemma 2.2, we can choose $M \gg 1$ such that, for $y \ge \delta$,

$$\frac{8|r'w'(y)|}{w(y)} \le M, \qquad w(y) \le My^k.$$

Then, for t_0 such that

$$1 + \frac{M}{\sqrt{t_0}} \le 2$$

we get, for any $t \ge t_0$ and any $z = x - 2t + r' \ln t \ge \delta \sqrt{t}$, that

$$\begin{split} \mathcal{L}\underline{u}(t,x) &= \varepsilon \frac{e^{-z}}{t^3} \left\{ -\frac{M}{2} w \left(\frac{z}{\sqrt{t}} \right) + \left(1 + \frac{M}{\sqrt{t}} \right) \left[r'w' \left(\frac{z}{\sqrt{t}} \right) + \left(1 + \frac{M}{\sqrt{t}} \right) \varepsilon e^{-z} w^2 \left(\frac{z}{\sqrt{t}} \right) \right] \right\} \\ &\leq \varepsilon \frac{e^{-z}}{t^3} w \left(\frac{z}{\sqrt{t}} \right) \left\{ -\frac{M}{4} + 4M \varepsilon e^{-z} \frac{z^k}{t^{k/2}} \right\}. \end{split}$$

Up to increasing t_0 if necessary, we have for any $t \ge t_0$ that $\max_{z\ge\delta\sqrt{t}} e^{-z} z^k = e^{-\delta\sqrt{t}} (\delta\sqrt{t})^k$, and then (up to also reducing $\varepsilon > 0$),

$$\mathcal{L}\underline{u}(t,x) \le \varepsilon \frac{e^{-z}}{t^3} w\left(\frac{z}{\sqrt{t}}\right) \left\{-\frac{M}{4} + 4M\varepsilon e^{-\delta\sqrt{t}}\delta^k\right\} < 0,$$

for any $x \ge 2t - r' \ln t + \delta \sqrt{t}$.

Putting the inequalities in both subdomains together, we found M > 0 such that, for any $\varepsilon > 0$ and $t_0 > 1$ respectively small and large enough, the function (4.6) is a sub-solution of (1.1) for $t \ge t_0$ and $x \ge 2t - r' \ln t$.

Step 2: comparison with the solution. We claim that, up to reducing $\varepsilon > 0$ and increasing $t_0 > 1$, we have that

$$u_0 \ge \underline{u}(t_0, \cdot), \tag{4.8}$$

where \underline{u} is defined in (4.6). Recall that $\underline{u}(t, x)$ is supported in $\{x \ge 2t - r' \ln t\}$, and hereafter we extend it continuously by 0 on $\{x < 2t - r' \ln t\}$. Therefore, by Assumption 1.1 and up to increasing t_0 , we only need to show that

$$ax^k e^{-x} \ge \underline{u}(t_0, x), \quad \forall x \ge 2t_0 - r' \ln t_0 \ge 1,$$

or equivalently,

$$ax^k \ge C(t_0)\varepsilon w\left(\frac{x-2t_0+r'\ln t_0}{\sqrt{t_0}}\right), \quad \forall x \ge 2t_0-r'\ln t_0 \ge 1,$$
(4.9)

where $C(t_0) := \left(1 + \frac{M}{\sqrt{t_0}}\right) e^{2t_0} t_0^{\frac{1}{2}-r}$. It follows from Lemma 2.2, more precisely (2.2), that (4.9) is satisfied if ε is small enough.

We conclude that (4.8) holds. Then, by the parabolic comparison principle,

$$u(t,x) \ge \underline{u}(t+t_0,x),$$

for any t > 0 and $x \in \mathbb{R}$. One can verify that $\|\underline{u}(t + t_0, \cdot)\|_{L^{\infty}(\mathbb{R})}$ decreases algebraically to 0 as $t \to +\infty$, due to r' = r - 2. In particular, we cannot yet conclude anything on the largetime position of the level sets. However, this decay in time, and specifically the choice of r', ensured that the nonlinear component of the reaction term was small enough when evaluating $\mathcal{L}\underline{u}$. Therefore, it played a crucial role in \underline{u} being a sub-solution.

Still, we now know that, for any t > 0 and $x = 2t + \sqrt{t}$,

$$u(t, 2t + \sqrt{t}) \ge \underline{u}(t + t_0, 2t + \sqrt{t}) \ge \varepsilon e^{-\sqrt{t}} e^{2t_0} (t + t_0)^{\frac{1}{2} - r} w\left(\frac{\sqrt{t} - 2t_0 + r' \ln(t + t_0)}{\sqrt{t + t_0}}\right)$$

and then

$$u(t, 2t + \sqrt{t}) \ge \frac{1}{2} \varepsilon e^{2t_0} w(1) e^{-\sqrt{t}} t^{\frac{1}{2} - r},$$

provided t is large enough. This concludes the proof of (4.3).

4.2 **Proof of Proposition 4.1**

We now address how the lower estimate on the level sets of the solution to (1.1) follows from Lemma 4.2. The main idea is to construct a second sub-solution on a left half-line, which turns out to be more challenging in the advance case.

Step 3: the second sub-solution in the delay case $0 \le r < \frac{3}{2}$. In this case, the minimal traveling wave itself serves as a sub-solution, see Lemma 2.1. Denote by U the solution to

$$\begin{cases} U'' + 2U' + U(1 - U) = 0 & \text{on } \mathbb{R}, \\ U(-\infty) = 1, \quad U(0) = \frac{1}{2}, \quad U(+\infty) = 0, \end{cases}$$

which is known to satisfy U' < 0 and $U(z) \sim Bze^{-z}$ as $z \to +\infty$, for some B > 0. Then, for any $0 < \alpha < 1$ to be selected below, and due to $r \ge 0$, it is straightforward to check that

$$\underline{u}_2(t,x) := \alpha U(x - 2t + r\ln(1+t))$$

is a sub-solution of (1.1) for $t \ge 0$ and $x \in \mathbb{R}$.

Next, we aim at applying the parabolic comparison principle on $\{(t, x) : t \ge T, x \le 2t + \sqrt{t}\}$ for some appropriate $T > T_0$, where T_0 is chosen as in Lemma 4.2. On the one hand, on the lateral boundary $x = 2t + \sqrt{t}$, we have

$$\underline{u}_2(t, 2t + \sqrt{t}) = \alpha U(\sqrt{t} + r \ln(1+t)) \sim \alpha B t^{\frac{1}{2} - r} e^{-\sqrt{t}}, \quad \text{ as } t \to +\infty.$$

In view of the lower estimate (4.3) of the solution u, it follows that, by setting $0 < \alpha < \frac{\varepsilon}{2B}$, we have, for some $T > T_0$ large enough, that

$$u(t, 2t + \sqrt{t}) \ge \underline{u}_2(t, 2t + \sqrt{t}), \quad \forall t \ge T$$

On the other hand, since spreading occurs (at speed c = 2) we have $\inf_{x \le 0} u(T, x) \ge \frac{1}{2}$, up to increasing T again. Such a time T now being fixed, and since $u(t, \cdot) > 0$ for any t > 0 by the strong maximum principle, we get that

$$\inf_{x \le 2T + \sqrt{T}} u(T, x) > 0$$

Hence, imposing $0 < \alpha < \inf_{x \leq 2T + \sqrt{T}} u(T, x)$, we immediately get that $u(T, x) \geq \underline{u}_2(T, x)$ for all $x \leq 2T + \sqrt{T}$.

As a result, for any $0 < \alpha < \min\{\frac{\varepsilon}{2B}, \inf_{x \leq 2T + \sqrt{T}} u(T, x)\}$, there exists T > 0 such that we can deduce from the comparison principle that

$$u(t,x) \ge \alpha U(x - 2t + r\ln(1+t)), \quad \forall t \ge T, \ \forall x \le 2t + \sqrt{t}.$$

$$(4.10)$$

Due to the asymptotics of U as $y \to +\infty$, there exists some $k_2 > 0$ small enough such that

$$U(y) \ge \frac{2k_2}{\alpha} y e^{-y}$$
 for any $y \ge 0$.

Thus, from (4.10), we get

$$u(t, 2t - r \ln t + y) \ge \alpha U(y + r \ln(1 + 1/t)) \ge k_2 y e^{-y},$$

for any t large enough and $0 \le y \le \sqrt{t} + r \ln t$. We infer that (4.2) holds.

Moreover, (4.1) is a direct consequence of (4.10) for "small" level sets, say $0 < m < m_0 := \frac{\alpha}{2}$, and then for all level sets 0 < m < 1 thanks to the following result (which does not require $r \ge 0$ and will be used again later).

Lemma 4.3. Let $0 < m_0 < 1$ be given. Then, if (4.1) holds for any $0 < m < m_0$, it actually holds for any 0 < m < 1.

Proof. Let us consider the remaining case $m_0 \le m < 1$. Let us denote by v = v(t, x) the solution to (1.1) starting from

$$v_0(x) := \begin{cases} \frac{m_0}{2} & \text{if } x \le -1, \\ -\frac{m_0}{2}x & \text{if } -1 < x < 0, \\ 0 & \text{if } x \ge 0. \end{cases}$$
(4.11)

Since spreading occurs (at speed c = 2), there is a time $\tau^* = \tau^*_{m_0,m} > 0$ such that

$$v(\tau^*, x) > m, \quad \forall x \le 0.$$

$$(4.12)$$

From (4.1) with $m = \frac{3}{4}m_0$, we know that there are $C_0 > 0$ and $T_0 > \max\{1, \tau^*\}$ such that

$$u(t,x) \ge \frac{3}{4}m_0, \quad \forall t \ge T_0, \ \forall x \le x^*(t) := 2t - r\ln t - C_0.$$

In view of definition (4.11), this implies

$$u(t,x) \ge v_0(x-x^*(t)), \quad \forall t \ge T_0, \ \forall x \in \mathbb{R},$$

so that the comparison principle yields

$$u(t+\tau, x) \ge v(\tau, x - x^*(t)), \quad \forall t \ge T_0, \ \forall \tau \ge 0, \ \forall x \in \mathbb{R}.$$

Using (4.12), we get that

$$u(t+\tau^*, x) > m, \quad \forall t \ge T_0, \ \forall x \le x^*(t),$$

which proves the result.

Step 4: the second sub-solution in the advance case r < 0. We look for a sub-solution that is valid for t > 0 and $x \le 2t + \sqrt{t}$. In the advance case, the traveling wave cannot be turned into an appropriate sub-solution. Instead, an adequate tool for the construction is the following ODE solution.

Lemma 4.4. Let $\gamma > 1$ be given. Denote $\varphi = \varphi_{\gamma}$ the solution of the ODE Cauchy problem

$$\begin{cases} \varphi'' + 2\varphi' + \varphi - \gamma \varphi^2 = 0 \quad on \ (0, +\infty), \\ \varphi(0) = \frac{1}{2\gamma}, \\ \varphi'(0) = 0. \end{cases}$$

Then

$$\varphi'(z) < 0, \qquad \frac{\varphi'(z)}{\varphi(z)} \ge -1, \quad \forall z > 0.$$

Moreover, there exists $B = B_{\gamma} > 0$ such that

$$\varphi(z) \sim Bze^{-z}, \quad as \ z \to +\infty.$$
 (4.13)

Proof. This is rather standard from the same ODE techniques which have been extensively used to characterize traveling waves of reaction-diffusion equations, but we will sketch the proof for the sake of completeness. First, the ODE is recast

$$\begin{cases} p' = q, \\ q' = -2q - p + \gamma p^2, \end{cases}$$

and we use a phase plane analysis. The trajectory corresponding to φ in the (p,q) plane starts from $(\frac{1}{2\gamma}, 0)$ and immediately enters the subset $\{0 . Then it cannot cross$ $the horizontal segment <math>\{0 , nor the vertical half-line <math>\{p = \frac{1}{\gamma}, q < 0\}$. It also cannot cross the diagonal $\{q = -p\}$ (if so then, at the first crossing point, we would have $-1 \leq \frac{q'}{p'} = \frac{-q + \gamma p^2}{q} = -1 + \gamma \frac{p^2}{q}$, which is impossible since q < 0). It already follows that $\varphi' < 0$ and $\varphi'/\varphi \geq -1$, and that $\varphi(+\infty) = 0$. By standard ODE perturbation theory, we also get that

$$\varphi(z) \sim (Bz+C)e^{-z}, \quad \text{as } z \to +\infty,$$

for some $B \ge 0$. We assume by contradiction that B = 0. Since the PDE $\partial_t u = \partial_{xx} u + u - \gamma u^2$ is still of the KPP type, it admits a minimal traveling wave $U = U_{\gamma}$ solving

$$\begin{cases} U'' + 2U' + U - \gamma U^2 = 0 \quad \text{on } \mathbb{R}, \\ U(-\infty) = \frac{1}{\gamma}, \quad U(0) = \frac{1}{2\gamma}, \quad U(+\infty) = 0, \end{cases}$$

which is known to satisfy U' < 0 and $U(z) \sim Aze^{-z}$ as $z \to +\infty$, for some A > 0. Hence φ and U solve the same ODE and, from the behaviors at z = 0 and $z \to +\infty$, the trajectories in the phase plane should intersect, which contradicts the Cauchy-Lipschitz theorem. We conclude that B > 0, and (4.13) holds.

Equipped with this, we now claim that we can find $\eta > 0$ and $\gamma > 1$ such that

$$\underline{u}_2(t,x) := \begin{cases} e^{\eta \frac{z}{\sqrt{t}}} \varphi(z) & \text{if } z > 0, \\ \frac{1}{2\gamma} & \text{if } z \le 0, \end{cases} \quad \text{where } z := x - 2t + r \ln t, \tag{4.14}$$

is a sub-solution of (1.1) for t > 0, $z \le 2\sqrt{t}$. The function \underline{u}_2 is clearly a sub-solution for z < 0, it is continuous and has the good angle at the gluing point in the sense that

$$\partial_x \underline{u}_2(t, (2t - r \ln t)^+) > 0 = \partial_x \underline{u}_2(t, (2t - r \ln t)^-), \quad \forall t > 0,$$

since $\varphi(0) = \frac{1}{2\gamma}$ and $\varphi'(0) = 0$. Thus we only need to deal with the differential inequality when $0 < z \le 2\sqrt{t}$. By using Lemma 4.4 and the ODE satisfied by φ , we compute

$$\begin{split} e^{-\eta \frac{z}{\sqrt{t}}} \mathcal{L}\underline{u}_{2}(t,x) &= e^{-\eta \frac{z}{\sqrt{t}}} \left(\partial_{t}\underline{u}_{2} - \partial_{xx}\underline{u}_{2} - \underline{u}_{2} + \underline{u}_{2}^{2}\right) \\ &= -\frac{\eta z}{2t^{3/2}}\varphi(z) - \frac{2\eta}{\sqrt{t}}\varphi'(z) - \frac{\eta^{2}}{t}\varphi(z) - \varphi''(z) \\ &- \left(2 - \frac{r}{t}\right) \left(\frac{\eta}{\sqrt{t}}\varphi(z) + \varphi'(z)\right) - \varphi(z) + e^{\eta \frac{z}{\sqrt{t}}}\varphi^{2}(z) \\ &= -\frac{\eta z}{2t^{3/2}}\varphi(z) - \frac{2\eta}{\sqrt{t}}(\varphi'(z) + \varphi(z)) - \frac{\eta^{2}}{t}\varphi(z) + \frac{r}{t}\varphi'(z) + \frac{r\eta}{t^{3/2}}\varphi(z) \\ &+ \left(e^{\eta \frac{z}{\sqrt{t}}} - \gamma\right)\varphi^{2}(z). \end{split}$$

Recalling that r < 0, z > 0, $\varphi > 0$, $\varphi' < 0$ and $\varphi' + \varphi \ge 0$, we get rid of the some nonpositive terms and find that

$$\frac{e^{-\eta \frac{z}{\sqrt{t}}}}{\varphi(z)} \mathcal{L}\underline{u}_2(t,x) \leq \frac{-\eta^2 - r}{t} + \left(e^{\eta \frac{z}{\sqrt{t}}} - \gamma\right)\varphi(z).$$

Picking $\eta = \sqrt{-r} > 0$ and $\gamma = e^{2\eta} > 1$, and using $z \le 2\sqrt{t}$, we get

$$\frac{e^{-\eta \frac{z}{\sqrt{t}}}}{\varphi(z)}\mathcal{L}\underline{u}_{2}(t,x) \leq (e^{2\eta}-\gamma)\varphi(z) = 0.$$

As a result, \underline{u}_2 is a sub-solution of (1.1) for t > 0, $x \le 2t - r \ln t + 2\sqrt{t}$.

Now, the conclusion is rather similar to that in the delay case. Precisely, we fix $\eta > 0$, $\gamma > 1$ as above. Recall that B > 0 is set from (4.13) and $\varepsilon > 0$ is set from Lemma 4.2. Let $0 < \alpha < 1$ be any small enough constant so that

$$\alpha e^{\eta}B < \frac{\varepsilon}{2}.$$

Observe first that $\alpha \underline{u}_2$ is also a sub-solution for $t \ge 0$, $x \le 2t - r \ln t + 2\sqrt{t}$. We aim at applying the parabolic comparison principle on $\{(t, x) : t \ge T, x \le 2t + \sqrt{t}\}$ for some appropriate $T > T_0$, where T_0 is defined as in Lemma 4.2. On the lateral boundary $x = 2t + \sqrt{t}$, we have

$$\alpha \underline{u}_2(t, 2t + \sqrt{t}) = \alpha e^{\eta (1 + r \frac{\ln t}{\sqrt{t}})} \varphi(\sqrt{t} + r \ln t) \sim \alpha e^{\eta} B t^{\frac{1}{2} - r} e^{-\sqrt{t}}, \quad \text{as } t \to +\infty.$$

In view of the lower estimate (4.3) of the solution u and our choice of α , we can find some $T > T_0$ large enough such that

$$u(t, 2t + \sqrt{t}) \ge \alpha \underline{u}_2(t, 2t + \sqrt{t}), \quad \forall t \ge T.$$

Moreover, since spreading occurs (at speed c = 2), up to increasing T (that we now fix), we have $\inf_{x \leq 0} u(T, x) \geq \frac{1}{2}$, and then $\inf_{x \leq 2T + \sqrt{T}} u(T, x) > 0$ thanks to the positivity of u. Hence,

by setting further $0 < \alpha < 2\gamma e^{-2\eta} \inf_{x \leq 2T + \sqrt{T}} u(T, x)$, one may check that $u(T, x) \geq \alpha \underline{u}_2(T, x)$ for all $x \leq 2T + \sqrt{T}$.

As a result, we deduce from the comparison principle that

$$u(t,x) \ge \alpha \underline{u}_2(t,x), \quad \forall t \ge T, \ \forall x \le 2t + \sqrt{t}.$$

$$(4.15)$$

The end of the proof proceeds as in the delay case. On the one hand, (4.1) immediately follows from (4.15) for "small" level sets, and then for all level sets thanks to Lemma 4.3. On the other hand, (4.2) follows from (4.15) together with the definition of \underline{u}_2 in (4.14) and the asymptotics of φ by Lemma 4.4. We omit the details.

5 The critical case k = -2

In this section, we prove the following proposition, from which Theorem 1.3 will follow.

Proposition 5.1 (Level sets in the critical case). Let u = u(t, x) be the solution to (1.1) starting from u_0 satisfying Assumption 1.1 with k = -2. Let $m \in (0, 1)$ be given. Then, there are $C_m > 0$ and $T_m > 0$ such that

$$E_m(t) \subset \left(2t - \frac{3}{2}\ln t + \ln\ln t - C_m, 2t - \frac{3}{2}\ln t + \ln\ln t + C_m\right), \quad \forall t \ge T_m.$$
(5.1)

Moreover, there are T > 0, $k_1 > k_2 > 0$ and $\sigma_1 > 0$ such that

$$k_2 y e^{-y} \le u \left(t, 2t - \frac{3}{2} \ln t + \ln \ln t + y \right) \le k_1 (y+1) e^{-y},$$
 (5.2)

for any $t \ge T$ and $0 \le y \le \sigma_1 \ln t$.

Step 1. Preliminary computations. Let us first consider v = v(t, x) the solution to the heat equation on the half-line with Dirichlet boundary condition, namely

$$\begin{cases} \partial_t v = \partial_{xx} v, & t > 0, \ x > 0, \\ v(t,0) = 0, & t > 0, \end{cases}$$
$$v_0(x) = \begin{cases} 1 & \text{if } 0 < x \le 1, \\ \frac{1}{x^2} & \text{if } x > 1. \end{cases}$$
(5.3)

starting from

In particular, the function
$$e^{-(x-2t)}v(t, x-2t)$$
 solves the linearized equation $\partial_t u = \partial_{xx}u + u$
on a right half-line with a moving Dirichlet condition at $x = 2t$, and its initial value is also
consistent with Assumption 1.1 with $k = -2$. Therefore, it serves as a potential candidate to
accurately determine the position of the level sets.

First, we know that

$$v(t,x) = \frac{1}{\sqrt{4\pi t}} \int_0^{+\infty} \left(e^{-\frac{(x-y)^2}{4t}} - e^{-\frac{(x+y)^2}{4t}} \right) v_0(y) dy$$
$$= \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi t}} \int_0^{+\infty} e^{-\frac{y^2}{4t}} \left(e^{\frac{yx}{2t}} - e^{-\frac{yx}{2t}} \right) v_0(y) dy,$$

so that

$$\begin{aligned} v(t,2\sqrt{t}) &= \frac{e^{-1}}{\sqrt{4\pi t}} \int_{0}^{+\infty} e^{-\frac{y^{2}}{4t}} \left(e^{\frac{y}{\sqrt{t}}} - e^{-\frac{y}{\sqrt{t}}} \right) v_{0}(y) dy \\ &= \frac{e^{-1}}{\sqrt{4\pi}} \int_{0}^{+\infty} e^{-\frac{y^{2}}{4}} \left(e^{y} - e^{-y} \right) v_{0}(\sqrt{t}y) dy \\ &= \frac{e^{-1}}{\sqrt{4\pi}} \int_{0}^{1/\sqrt{t}} e^{-\frac{y^{2}}{4}} \left(e^{y} - e^{-y} \right) dy + \frac{e^{-1}}{\sqrt{4\pi}} \int_{1/\sqrt{t}}^{+\infty} e^{-\frac{y^{2}}{4}} \left(e^{y} - e^{-y} \right) \frac{1}{ty^{2}} dy \\ &=: I_{1} + I_{2}. \end{aligned}$$

On the one hand, as $t \to +\infty$,

$$I_1 \sim \frac{e^{-1}}{\sqrt{\pi}} \int_0^{1/\sqrt{t}} y dy = \mathcal{O}\left(\frac{1}{t}\right).$$

On the other hand, due to the integrand in I_2 being integrable at $+\infty$, but not at 0, we have, as $t \to +\infty$,

$$I_2 \sim \frac{e^{-1}}{\sqrt{4\pi}} \int_{1/\sqrt{t}}^1 e^{-\frac{y^2}{4}} \left(e^y - e^{-y} \right) \frac{1}{ty^2} dy \sim \frac{e^{-1}}{\sqrt{\pi}} \int_{1/\sqrt{t}}^1 \frac{1}{ty} dy \sim \frac{e^{-1}}{\sqrt{\pi}} \frac{\ln t}{2t}.$$

Therefore

$$v(t, 2\sqrt{t}) \sim \frac{e^{-1}}{\sqrt{\pi}} \frac{\ln t}{2t}, \quad \text{as } t \to +\infty.$$
 (5.4)

Let us conclude this step by a formal discussion. If the ansatz $e^{-(x-2t)}v(t, x-2t)$ is accurate, we may expect $u(t, 2t + 2\sqrt{t})$ to be well approximated by

$$e^{-2\sqrt{t}}v(t,2\sqrt{t}) \sim \frac{e^{-1}}{\sqrt{\pi}}e^{-2\sqrt{t}}\frac{\ln t}{2t}, \quad \text{as } t \to +\infty.$$

On the other hand, we have that

$$U\left(2t + 2\sqrt{t} - \left(2t - \frac{3}{2}\ln t + \ln\ln t\right)\right) = U\left(2\sqrt{t} + \frac{3}{2}\ln t - \ln\ln t\right)$$
$$\sim B \times \left(2\sqrt{t} + \frac{3}{2}\ln t - \ln\ln t\right)e^{-2\sqrt{t} - \frac{3}{2}\ln t + \ln\ln t}$$
$$\sim 2Be^{-2\sqrt{t}}\frac{\ln t}{t}, \quad \text{as } t \to +\infty.$$

Observing that both estimates coincide, this suggests that the position of the front of the solution u should be around $2t - \frac{3}{2} \ln t + \ln \ln t$, in agreement with our Theorem 1.3.

Step 2. Lower estimate. The rough plan is to first perturb the function v defined in the previous step to make it a sub-solution, thereby obtaining a lower bound on the solution at $x = 2t + 2\sqrt{t}$. Subsequently, the lower estimate on the level sets will follow similarly to the noncritical delay case, thanks to the fact that a slowed-down minimal front also serves as a sub-solution.

Precisely, we consider

$$\underline{u}(t,x) = \delta \times \left(1 + \frac{M}{(t+1)^{1/4}}\right) \times e^{-(x-2t)}v(t,x-2t), \quad t \ge 0, x \ge 2t,$$

where $\delta > 0$ and M > 0 will be specified below, chosen to be respectively small and large enough. Notice that we multiplied the ansatz $e^{-(x-2t)}v(t, x - 2t)$ by a decreasing in time function. As in the previous section, provided that v decays to 0 fast enough, the resulting additional term when evaluating the KPP equation will absorb the nonlinear quadratic term that our computation in Step 1 had previously ignored.

First, by Assumption 1.1 (with k = -2) and our choice of $v(0, \cdot)$, it is straightforward to check that, for any M > 0, there exists $\delta > 0$ small enough so that

$$\underline{u}(0,x) \le u_0(x), \quad \forall x \ge 0.$$

By the strong maximum principle, we also have that

$$\underline{u}(t,2t) = 0 < u(t,2t), \quad \forall t > 0.$$

Next, let us show that \underline{u} is a sub-solution. We compute, for any t > 0 and x > 2t,

$$\begin{split} \delta^{-1} e^{x-2t} &\times \left(\partial_t \underline{u} - \partial_{xx} \underline{u} - \underline{u} + \underline{u}^2\right) \\ &= v(t, x - 2t) \times \left[\delta \left(1 + \frac{M}{(t+1)^{1/4}}\right)^2 e^{-(x-2t)} v(t, x - 2t) - \frac{M}{4(t+1)^{5/4}}\right] \\ &\leq v(t, x - 2t) \times \left[e^{-(x-2t)} v(t, x - 2t) - \frac{M}{4(t+1)^{5/4}}\right], \end{split}$$

assuming further $\delta(1+M)^2 < 1$. From the following lemma, whose proof is postponed, we can choose M > 0 sufficiently large (and then $\delta > 0$ sufficiently small) to ensure that \underline{u} is a sub-solution.

Lemma 5.2. For any $\varepsilon > 0$, the function v defined in Step 1 satisfies

$$\sup_{t>0, x>0} e^{-x} v(t,x) \times (t+1)^{\frac{3}{2}-\varepsilon} < +\infty.$$

Hence, by the parabolic comparison principle, we have that

$$u(t,x) \ge \underline{u}(t,x) = \delta\left(1 + \frac{M}{(t+1)^{1/4}}\right) e^{-(x-2t)} v(t,x-2t), \text{ for all } t \ge 0, x \ge 2t$$

As a result, we deduce from (5.4) that that there exists $\varepsilon > 0$ such that

$$u(t, 2t + 2\sqrt{t}) \ge \varepsilon e^{-2\sqrt{t}} \times \frac{\ln t}{t}, \quad \forall t > 1.$$
(5.5)

The rest of the proof proceeds similarly to the noncritical case, see Step 3 of subsection 4.2 for further details. Consider the second sub-solution (as can be easily verified)

$$\underline{u}_2(t,x) := \alpha U\left(x - 2t + \frac{3}{2}\ln t - \ln\ln t\right)$$

where $\alpha \in (0,1)$ and U denotes the minimal travelling wave. Recalling $U(z) \sim Bze^{-z}$ as $z \to +\infty$, we have

$$\underline{u}_2(t, 2t + 2\sqrt{t}) \sim 2\alpha B e^{-2\sqrt{t}} \times \frac{\ln t}{t}, \quad \text{as } t \to +\infty.$$

In view of (5.5), we may thus choose α small enough so that, for some large enough T > 0,

$$\underline{u}_2(t, 2t + 2\sqrt{t}) \le u(t, 2t + 2\sqrt{t}), \quad \forall t \ge T.$$

Up to further reducing α , we may also have

$$\underline{u}_2(T,x) \le u(T,x),$$

for $x \leq 2T + 2\sqrt{T}$. Finally, by another application of the parabolic comparison principle, we have that

$$\underline{u}_2(t,x) \le u(t,x), \quad \text{for any } t \ge T \text{ and } x \le 2t + 2\sqrt{t}.$$
(5.6)

We can now conclude that there exist $C_m, T_m > 0$

$$E_m(t) \subset \left(2t - \frac{3}{2}\ln t + \ln\ln t - C_m, +\infty\right), \text{ for all } t \ge T_m,$$

first for small $0 < m < \frac{\alpha}{2}$, then for all 0 < m < 1 by Lemma 4.3. Let us point out that this already proves the drift is not exactly logarithmic.

It also follows from (5.6) and the spatial asymptotics of U (hence, of \underline{u}_2) that there exists $k_2 > 0$ such that

$$u\left(t, 2t - \frac{3}{2}\ln t + \ln\ln t + y\right) \ge k_2 y e^{-y},$$

for all $t \ge T$ and $0 \le y \le \sqrt{t}$. In particular, we have proved the first inequality in (5.2).

Step 3. Upper estimate. Let us consider

$$\widehat{u}(t,x) := \left(1 - \frac{M}{t^{1/4}}\right) \times \frac{t^{3/2}}{\ln t} e^{-(x - 2t + \frac{3}{2}\ln t - \ln\ln t)} v\left(t, x - 2t + \frac{3}{2}\ln t - \ln\ln t\right),$$

where M > 0 is sufficiently large (to be specified later), and show that it is a positive supersolution for t large enough and $x > 2t - \frac{3}{2} \ln t + \ln \ln t$. Denoting $z = x - 2t + \frac{3}{2} \ln t - \ln \ln t$ for convenience, it is straightforward to compute

$$\left(1 - \frac{M}{t^{1/4}}\right)^{-1} \times \frac{e^{z} \ln t}{t^{3/2}} \times \left(\partial_{t}\widehat{u} - \partial_{xx}\widehat{u} - \widehat{u}\right) = \left(\frac{3}{2t} - \frac{1}{t\ln t}\right)\partial_{x}v + \frac{M}{4t^{5/4}}\left(1 - \frac{M}{t^{1/4}}\right)^{-1}v(t, z),$$

for t > 0 and z > 0. From the following lemma, whose proof is postponed, we may choose $t_0 > 0$ and M > 0 sufficiently large to make \hat{u} a (positive) super-solution of the KPP equation on the moving right half-line $t > t_0$, $x > 2t - \frac{3}{2} \ln t + \ln \ln t$.

Lemma 5.3. There are $t_0 > 0$ and C > 0 such that the function v defined in Step 1 satisfies

$$\frac{\partial_x v(t,x)}{v(t,x)} \ge -\frac{C}{t^{1/4}}, \quad \text{ for all } t > t_0, \ x > 0.$$

We now sketch the end of the proof of the upper estimate. First, for any K > 1, we define

$$\overline{u}(t,x) = \begin{cases} 1 & \text{if } x < 2(t+t_0) - \frac{3}{2}\ln(t+t_0) + \ln\ln(t+t_0) + 1, \\ \min\{1, K\widehat{u}(t+t_0, x)\} & \text{if } x \ge 2(t+t_0) - \frac{3}{2}\ln(t+t_0) + \ln\ln(t+t_0) + 1. \end{cases}$$

We have that

$$K\widehat{u}\left(t+t_0, 2(t+t_0) - \frac{3}{2}\ln(t+t_0) + \ln\ln(t+t_0) + 1\right)$$
$$= K\left(1 - \frac{M}{(t+t_0)^{1/4}}\right) \times \frac{(t+t_0)^{3/2}}{\ln(t+t_0)}e^{-1}v(t+t_0, 1)$$

which, up to increasing t_0 and K, is strictly larger than 1 for all t > 0 thanks to the following lemma, whose proof is also postponed. As a result, \overline{u} is a generalized super-solution of the KPP equation (for all t > 0 and $x \in \mathbb{R}$), for any choice of K large enough.

Lemma 5.4. There are $C_2 > C_1 > 0$ and T > e such that the function v defined in Step 1 satisfies, for any $t \ge T$ and $x \in (1, \ln t)$,

$$C_1 \frac{x \ln t}{t^{3/2}} \le v(t, x) \le C_2 \frac{x \ln t}{t^{3/2}}.$$

In order to apply the comparison principle, let us check that

$$u_0 \le \overline{u}(t_0, \cdot) \quad \text{on } \mathbb{R}.$$
 (5.7)

Proceeding as in the noncritical case, up to increasing the constant K, it suffices to check this inequality in a neighborhood of $+\infty$. With the time t_0 now fixed, by the strong maximum principle and the positivity of the initial data v_0 , see (5.3), there exists $\alpha \in (0, 1)$ small enough such that

$$v(t,1) \ge \alpha, \quad \forall t \in [0,t_0].$$

Moreover, recall that

$$v_0(x) = \frac{1}{x^2}, \quad \forall x \ge 1.$$

Also, since $\underline{v}(t, x) := \frac{\alpha}{x^2}$ satisfies

$$\partial_t \underline{v} = 0 < \partial_{xx} \underline{v}, \quad x \ge 1,$$

by the comparison principle, we get that

$$v(t_0, x) \ge \frac{\alpha}{x^2}, \quad \forall x \ge 1.$$

In particular, for all $x \ge 2t_0 - \frac{3}{2} \ln t_0 + \ln \ln t_0 + 1$,

$$\widehat{u}(t_0, x) \ge \left(1 - \frac{M}{t_0^{1/4}}\right) \times \frac{t_0^{3/2}}{\ln t_0} e^{2t_0 - \frac{3}{2}\ln t_0 + \ln\ln t_0} \times \frac{\alpha e^{-x}}{(x - 2t_0 + \frac{3}{2}\ln t_0 - \ln\ln t_0)^2}.$$

We can thus find K large enough so that

$$K\widehat{u}(t_0, x) \ge A \frac{e^{-x}}{x^2} \ge u_0(x)$$

on a neighborhood of $+\infty$, so that (5.7) does hold.

We are now in a position to apply the comparison principle, yielding

$$u(t,x) \le \overline{u}(t+t_0,x),$$

for all t > 0 and $x \in \mathbb{R}$. The upper estimate on the position of the level sets finally follows, as detailed below.

First, from our results in the noncritical case (say for a chosen k in (-2, -1) for which r > 1) and the comparison principle, we already know that

$$\limsup_{t \to +\infty} \sup_{x \ge 2(t+t_0) - \ln(t+t_0)} u(t, x) = 0.$$
(5.8)

Next, for any $m \in (0, 1)$, select X > 1 large enough so that $KC_2Xe^{-X} < m$ (where C_2 comes from Lemma 5.4) and define the interval

$$I_t := \left(2t - \frac{3}{2}\ln t + \ln\ln t + X, 2t - \ln t\right).$$

Then we have

where we have used Lemma 5.4 again for the antepenultimate inequality. Putting this together with (5.8), we infer that there is $T_m > 0$ large enough so that, for all $t \ge T_m$,

$$E_m(t) \subset \left(-\infty, 2(t+t_0) - \frac{3}{2}\ln(t+t_0) + \ln\ln(t+t_0) + X\right),$$

which completes the proof of (5.1).

Now fixing, say, $m = \frac{1}{2}$ and proceeding similarly, we may also find that, for any t > 0 and y such that

$$2t - \frac{3}{2}\ln t + \ln\ln t + y \in I_{t+t_0},$$

then

for some $k_1 > 0$, and for all $t \ge T$ with T > 0 large enough. Now, letting $\sigma_1 \in (0, 1/2)$, we can assume up to increasing T that, for all $t \ge T$,

$$y \in (2t_0 + X + 1, \sigma_1 \ln t) \Rightarrow 2t - \frac{3}{2} \ln t + \ln \ln t + y \in I_{t+t_0},$$

and then

$$u\left(t, 2t - \frac{3}{2}\ln t + \ln\ln t + y\right) \le k_1(y+1)e^{-y},\tag{5.9}$$

for all $t \ge T$ and $2t_0 + X + 1 \le y \le \sigma_1 \ln t$. Since the left-hand term of (5.9) is bounded from above by 1, and the right-hand term has a positive minimum on $[0, 2t_0 + X + 1]$, one can deduce that (5.9) holds true for all $t \ge T$ and $0 \le y \le \sigma_1 \ln t$, up to increasing k_1 . This proves the second inequality in (5.2).

It thus only remains to prove the three postponed lemmas on the function v defined in Step 1.

Proof of Lemma 5.2. Notice that $e^{-(x-2t)}v(t, x-2t)$ solves the linearized equation

$$\partial_t u = \partial_{xx} u + u,$$

together with a Dirichlet boundary condition at x = 2t. Recall also that we have found in the noncritical case a family of super-solutions to this linear equation. Indeed, for any $k \in (-2, 1)$ and $r = \frac{1-k}{2} > 0$ (see Step 1 in subsection 3.1 in the delay case), then

$$\overline{u}(t,x) = \left(1 - \frac{M}{\sqrt{t}}\right) e^{-(x-2t+r\ln t)}\sqrt{t} \times w\left(\frac{x-2t+r\ln t}{\sqrt{t}}\right),$$

where w comes from Lemma 2.2, is a positive super-solution to this linear equation for all $t \ge t_0$ (with t_0 large enough) and $x \ge 2t > 2t - r \ln t$.

By construction,

$$e^{-x}v(0,x) \le x^k e^{-x},$$

for all x > 1. Therefore, proceeding similarly to the noncritical case, we may find that

$$e^{-(x-2t)}v(t,x-2t) \le K\overline{u}(t+t_0,x),$$

for some large enough K, and any t > 0, x > 2t. Due to the properties of w from Lemma 2.2, namely the fact that $\sup_{y \ge 0} w(y)/y < +\infty$, we get that

$$e^{-(x-2t)}v(t,x-2t) \leq Ke^{-(x-2(t+t_0)+r\ln(t+t_0))}\sqrt{t} \times w\left(\frac{x-2(t+t_0)+r\ln(t+t_0)}{\sqrt{t+t_0}}\right)$$

$$\leq Ce^{-(x-2t)} \times t^{-r} \times (x-2t+r\ln(t+t_0)),$$

for some C > 0. Since r can be chosen arbitrarily close to 3/2, we reach the conclusion. \Box *Proof of Lemma 5.3.* First, recall that

$$v(t,x) = \frac{1}{\sqrt{4\pi t}} \int_0^{+\infty} \left(e^{-\frac{(x-y)^2}{4t}} - e^{-\frac{(x+y)^2}{4t}} \right) v_0(y) dy$$
$$= \frac{e^{-\frac{x^2}{4t}}}{\sqrt{\pi t}} \int_0^{+\infty} e^{-\frac{y^2}{4t}} \sinh\left(\frac{xy}{2t}\right) v_0(y) dy.$$

Using the second expression, we find that

$$\partial_x v(t,x) = -\frac{x}{2t}v(t,x) + \frac{e^{-\frac{x^2}{4t}}}{\sqrt{\pi t}} \int_0^{+\infty} \frac{ye^{-\frac{y^2}{4t}}}{2t} \cosh\left(\frac{xy}{2t}\right) v_0(y) dy \ge -\frac{x}{2t}v(t,x).$$

It already follows that

$$\frac{\partial_x v(t,x)}{v(t,x)} \ge -\frac{1}{t^{1/4}} \quad \text{ for all } t > 0, \ 0 \le x \le 2t^{3/4}.$$

Next, let us consider the case $x > 2t^{3/4}$. On the one hand, for any t > 1 (so that also $x - \sqrt{t} > 1$), we have

$$\begin{aligned} v(t,x) &\geq \frac{1}{\sqrt{4\pi t}} \int_{x-\sqrt{t}}^{x+\sqrt{t}} \left(e^{-\frac{(x-y)^2}{4t}} - e^{-\frac{(x+y)^2}{4t}} \right) v_0(y) dy \\ &= \frac{1}{\sqrt{4\pi t}} \int_{-\sqrt{t}}^{\sqrt{t}} \left(e^{-\frac{y^2}{4t}} - e^{-\frac{(2x+y)^2}{4t}} \right) \times \frac{1}{(x+y)^2} dy \\ &\geq \frac{e^{-\frac{1}{4}}}{\sqrt{4\pi t}} \int_{-\sqrt{t}}^{\sqrt{t}} \left(1 - e^{-\frac{x^2+xy}{t}} \right) \times \frac{1}{(x+y)^2} dy. \end{aligned}$$

But $-\sqrt{t} \le y \le \sqrt{t}$, $x > 2t^{3/4}$ and t > 1 imply $\frac{x^2 + xy}{t} \ge 2t^{1/4} \ge 2$ so that

$$v(t,x) \ge \frac{e^{-\frac{1}{4}}}{x^2\sqrt{4\pi t}}(1-e^{-2})\int_{-\sqrt{t}}^{\sqrt{t}} \frac{1}{(1+\frac{y}{x})^2} dy \ge \frac{\delta}{x^2},$$

for some $\delta > 0$. On the other hand,

$$\begin{aligned} \partial_x v(t,x) &= \frac{1}{\sqrt{4\pi t}} \int_0^{+\infty} \left(-\frac{x-y}{2t} e^{-\frac{(x-y)^2}{4t}} + \frac{x+y}{2t} e^{-\frac{(x+y)^2}{4t}} \right) v_0(y) dy \\ &\geq -\frac{1}{\sqrt{4\pi t}} \int_0^{+\infty} \frac{x-y}{2t} e^{-\frac{(x-y)^2}{4t}} v_0(y) dy \\ &= -\frac{1}{\sqrt{4\pi t}} (I_1 + I_2 + I_3), \end{aligned}$$

where I_1 , I_2 , and I_3 are obtained by integrating, respectively, over (0, 1), $(x(1 - t^{-1/8}), x(1 + t^{-1/8}))$, and the remaining intervals. Recalling (5.3), we have, for any $t \ge t_0$ with $t_0 > 0$ large enough and $x > 2t^{3/4}$,

$$I_1 = \int_0^1 \frac{x - y}{2t} e^{-\frac{(x - y)^2}{4t}} dy \le \frac{x e^{-\frac{(x - 1)^2}{4t}}}{2t},$$

but also (in the sequel C denotes a positive constant that may vary from one line to another)

$$\begin{split} I_2 &= \int_{x(1-t^{-1/8})}^{x(1+t^{-1/8})} \frac{x-y}{2t} e^{-\frac{(x-y)^2}{4t}} v_0(y) dy \\ &= \int_{-xt^{-1/8}}^{xt^{-1/8}} \frac{y}{2t} e^{-\frac{y^2}{4t}} \times \frac{1}{x^2(1-\frac{y}{x})^2} dy \\ &\leq \int_{-xt^{-1/8}}^{xt^{-1/8}} \frac{y}{2t} e^{-\frac{y^2}{4t}} \times \frac{1+2\frac{y}{x}}{x^2} dy + C \int_{-xt^{-1/8}}^{xt^{-1/8}} \frac{|y|^3}{tx^4} e^{-\frac{y^2}{4t}} dy \\ &= 2 \int_0^{xt^{-1/8}} \frac{y^2}{tx^3} e^{-\frac{y^2}{4t}} dy + C \int_{-xt^{-1/8}}^{xt^{-1/8}} \frac{|y|^3}{tx^4} e^{-\frac{y^2}{4t}} dy \\ &\leq \frac{2\sqrt{t}}{x^3} \int_0^{+\infty} y^2 e^{-\frac{y^2}{4}} dy + \frac{Ct}{x^4} \int_{-\infty}^{+\infty} |y|^3 e^{-\frac{y^2}{4}} dy \\ &\leq \frac{C\sqrt{t}}{x^3} \quad (\text{since } x > 2t^{3/4}), \end{split}$$

and finally

$$\begin{split} I_{3} &= \left(\int_{1}^{x(1-t^{-1/8})} + \int_{x(1+t^{-1/8})}^{+\infty}\right) \frac{x-y}{2t} e^{-\frac{(x-y)^{2}}{4t}} v_{0}(y) dy \\ &\leq \left(\int_{1}^{x(1-t^{-1/8})} + \int_{x(1+t^{-1/8})}^{+\infty}\right) \frac{|x-y|}{2t} e^{-\frac{(x-y)^{2}}{4t}} dy \\ &\leq \left(\int_{-\infty}^{-xt^{-1/8}} + \int_{xt^{-1/8}}^{+\infty}\right) \frac{|y|}{2t} e^{-\frac{y^{2}}{4t}} dy \\ &\leq \left(\int_{-\infty}^{-xt^{-5/8}} + \int_{xt^{-5/8}}^{+\infty}\right) \frac{|z|}{2} e^{-\frac{z^{2}}{4}} dz \\ &\leq C e^{-\frac{(xt^{-5/8})^{2}}{4}}. \end{split}$$

Using again the fact that $x > 2t^{3/4}$, we have, for any $t \ge t_0$ with $t_0 > 0$ large enough,

$$\frac{x^3}{\sqrt{t}}I_1 \le \frac{x^4 e^{-\frac{(x-1)^2}{4t}}}{2t^{3/2}} \le C t^{3/2} e^{-\sqrt{t}} \le C,$$

and

$$\frac{x^3}{\sqrt{t}}I_3 \le C\frac{x^3}{\sqrt{t}}e^{-\frac{(xt^{-5/8})^2}{4}} \le Ct^{7/4}e^{-t^{1/4}} \le C.$$

Finally, putting all this together, we find that

$$\frac{\partial_x v(t,x)}{v(t,x)} \geq \frac{-(I_1 + I_2 + I_3)}{\sqrt{4\pi t} \times v(t,x)} \geq -\frac{C}{x^3} \times \frac{x^2}{\delta} \geq -\frac{C}{t^{3/4}},$$

for any $t \ge t_0$ with $t_0 > 0$ large enough and $x > 2t^{3/4}$. Putting this together with our previous estimate for $x \le 2t^{3/4}$, the proof of Lemma 5.3 is complete.

Proof of Lemma 5.4. Consider t > e and $x \in (1, \ln t)$. Then, we have

$$v(t,x) = \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi t}} \int_0^{+\infty} e^{-\frac{y^2}{4t}} \left(e^{\frac{yx}{2t}} - e^{-\frac{yx}{2t}} \right) v_0(y) dy$$

$$= \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi}} \int_0^{+\infty} e^{-\frac{y^2}{4}} \left(e^{\frac{yx}{2\sqrt{t}}} - e^{-\frac{yx}{2\sqrt{t}}} \right) v_0(\sqrt{t}y) dy$$

$$= I_1 + I_2 + I_3 + I_4,$$

where I_i $(1 \le i \le 4)$ will be defined and estimated below. For the sake of expediency, in the sequel $f(t, x) \approx g(t, x)$ means that there is C > 0 such that

 $\frac{f(t,x)}{Cg(t,x)} \to 1 \quad \text{as } t \to +\infty, \text{ uniformly with respect to } x \in (1, \ln t).$

Now, recalling (5.3),

$$\begin{split} I_1 &:= \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi}} \int_0^{1/\sqrt{t}} e^{-\frac{y^2}{4}} \left(e^{\frac{yx}{2\sqrt{t}}} - e^{-\frac{yx}{2\sqrt{t}}} \right) dy \\ &\approx \int_0^{1/\sqrt{t}} \frac{yx}{\sqrt{t}} dy \\ &\approx \frac{x}{t^{3/2}}, \end{split}$$

$$I_2 &:= \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi}} \int_{1/\sqrt{t}}^1 e^{-\frac{y^2}{4}} \left(e^{\frac{yx}{2\sqrt{t}}} - e^{-\frac{yx}{2\sqrt{t}}} \right) \times \frac{1}{ty^2} dy \\ &\approx \int_{1/\sqrt{t}}^1 \frac{x}{t^{3/2}y} dy \\ &\approx \frac{x \ln t}{t^{3/2}}, \end{split}$$

$$\leq I_3 &:= \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi}} \int_1^{\sqrt{t}/x} e^{-\frac{y^2}{4}} \left(e^{\frac{yx}{2\sqrt{t}}} - e^{-\frac{yx}{2\sqrt{t}}} \right) \times \frac{1}{ty^2} dy$$

$$0 \le I_3 := \frac{e^{-x}}{\sqrt{4\pi}} \int_1^{1} e^{-\frac{y}{4}} \left(e^{\frac{y}{2\sqrt{t}}} - e^{-\frac{y}{2\sqrt{t}}} \right) \times \frac{1}{ty^2} dy$$
$$\le C \int_1^{\sqrt{t}/x} e^{-\frac{y^2}{4}} \times \frac{x}{t^{3/2}y} dy \approx \frac{x}{t^{3/2}},$$

and, recalling also that $x \in (1, \ln t)$,

$$\begin{aligned} 0 &\leq I_4 &:= \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi}} \int_{\sqrt{t}/x}^{+\infty} e^{-\frac{y^2}{4}} \left(e^{\frac{yx}{2\sqrt{t}}} - e^{-\frac{yx}{2\sqrt{t}}} \right) \times \frac{1}{ty^2} dy \\ &\leq \frac{1}{t} \int_{\sqrt{t}/x}^{+\infty} e^{-\frac{y^2}{8}} dy \\ &\leq \frac{1}{t} \frac{x}{\sqrt{t}} \int_{\sqrt{t}/x}^{+\infty} y e^{-\frac{y^2}{8}} dy \\ &\leq C \frac{x}{t^{3/2}} e^{-\frac{t}{8\ln^2 t}}. \end{aligned}$$

As a result $v(t,x) = I_1 + I_2 + I_3 + I_4 \approx \frac{x \ln t}{t^{3/2}}$, which ends the proof.

6 Convergence of the profile

In this section, we complete the proof of our main results, namely Theorems 1.2 and 1.3. In both the noncritical and critical cases, we have already located the level sets, and derived some lower and upper estimates on the solution in the appropriate moving frame. Now, we address the convergence to a minimal traveling wave in that moving frame.

The proof, which relies on a Liouville-type result from [5], closely parallels the one in [19, Section 4] to which we refer for further details. Notably, it proceeds exactly in the same manner for both the noncritical and critical cases. Therefore, we only provide the details in the former, and point out that in the latter, one should use Proposition 5.1 instead of Propositions 3.1 and 4.1, and (5.2) instead of (3.2) and (4.2).

Proof of Theorem 1.2. First, recall from Propositions 3.1 and 4.1 that, for any 0 < m < 1, the *m*-level set $E_m(t)$ of the solution satisfies

$$E_m(t) \subset (2t - r\ln t - C_m, 2t - r\ln t + C_m), \quad \forall t \ge T_m,$$

$$(6.1)$$

for some $C_m > 0$ and $T_m > 0$.

We first set C > 0 large enough such that

$$Be^{-C} \le k_2 \le k_1 \le Be^C, \tag{6.2}$$

where k_1 and k_2 are defined in Propositions 3.1 and 4.1, and *B* comes from the asymptotics $U(z) \sim Bze^{-z}$ as $z \to +\infty$. We will show that (1.6) holds with this choice of *C*.

We argue by contradiction, and assume there are $\varepsilon > 0$ and a sequence of times $(t_n)_{n \in \mathbb{N}}$ such that $t_n \to +\infty$ as $n \to +\infty$ and

$$\inf_{|h| \le C} \left\| u(t_n, \cdot) - U(\cdot - 2t_n + r \ln t_n + h) \right\|_{L^{\infty}(0, +\infty)} \ge \varepsilon, \quad \forall n \in \mathbb{N}.$$

Since $U(-\infty) = 1$ and $U(+\infty) = 0$, by estimate (6.1) on the position of the level sets, we can find $\kappa > 0$ such that

$$\inf_{|k| \le C} \left(\max_{|x| \le \kappa} \left| u(t_n, x + 2t_n - r \ln t_n) - U(x+h) \right| \right) \ge \varepsilon, \quad \forall n \in \mathbb{N}.$$
(6.3)

Up to extraction, the sequence of functions

$$u_n(t,x) := u(t+t_n, x+2t_n - r\ln t_n), \quad t \in \mathbb{R}, x \in \mathbb{R},$$

converges locally uniformly in \mathbb{R}^2 to some $0 \leq u_{\infty} = u_{\infty}(t, x) \leq 1$ which solves (1.1) on \mathbb{R}^2 . Furthermore, we have

$$\lim_{y \to +\infty} \left(\sup_{(t,x) \in \mathbb{R}^2, \ x \ge 2t+y} u_{\infty}(t,x) \right) = 0 \quad \text{and} \quad \lim_{y \to -\infty} \left(\inf_{(t,x) \in \mathbb{R}^2, \ x \le 2t+y} u_{\infty}(t,x) \right) = 1.$$
(6.4)

Now, we fix $t \in \mathbb{R}$ and y > 0, and define

$$y_n := y + r \ln \frac{t + t_n}{t_n}.$$

For n sufficiently large, we have $t + t_n \ge 1$ and $0 \le y_n \le \min\{\sigma_1, \sigma_2\}\sqrt{t + t_n}$, so that we deduce from (3.2) and (4.2) that

$$k_2 y_n e^{-y_n} \le u_n(t, 2t+y) \le k_1(y_n+1)e^{-y_n}.$$

Passing to the limit as $n \to +\infty$, we obtain

$$k_2 y e^{-y} \le u_{\infty}(t, 2t+y) \le k_1(y+1)e^{-y}$$
 for any $t \in \mathbb{R}, y \ge 0.$ (6.5)

Therefore, a Liouville-type theorem, see [19, Lemma 4.1] and [5, Theorem 3.5], implies that there exists $h_0 \in \mathbb{R}$ such that the time-global solution u_{∞} satisfies

$$u_{\infty}(t,x) = U(x - 2t + h_0), \quad \forall (t,x) \in \mathbb{R}^2.$$
 (6.6)

We are now ready to complete the proof. Recall that $U(z) \sim Bze^{-z}$ as $z \to +\infty$. Thus, from (6.5) and (6.6), we have $k_2 \leq Be^{-h_0} \leq k_1$ which, in view of (6.2), enforces $|h_0| \leq C$. However, $u_n(0, x) \to u_\infty(0, x)$ uniformly for $x \in [-\kappa, \kappa]$ means nothing else than

$$\max_{|x| \le \kappa} \left| u(t_n, x + 2t_n - r \ln t_n) - U(x + h_0) \right| \to 0 \quad \text{as} \quad n \to +\infty,$$

which contradicts (6.3) since $|h_0| \leq C$. Therefore, the proof of Theorem 1.2 is complete. \Box

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References

- M. ABRAMOWITZ AND I. A. STEGUN, Handbook of mathematical functions with formulas, graphs, and mathematical tables, vol. 55 of National Bureau of Standards Applied Mathematics Series, For sale by the Superintendent of Documents, U.S. Government Printing Office, Washington, D.C., 1964.
- [2] J. AN, C. HENDERSON, AND L. RYZHIK, Front location determines convergence rate to traveling waves, 2023.
- [3] —, Quantitative steepness, semi-FKPP reactions, and pushmi-pullyu fronts, Arch. Ration. Mech. Anal., 247 (2023), p. 84. Id/No 88.
- [4] D. G. ARONSON AND H. F. WEINBERGER, Multidimensional nonlinear diffusion arising in population genetics, Adv. in Math., 30 (1978), pp. 33–76.
- [5] H. BERESTYCKI AND F. HAMEL, Generalized travelling waves for reaction-diffusion equations, in Perspectives in nonlinear partial differential equations in honor of Haïm Brezis. Based on the conference celebration of Haïm Brezis' 60th birthday, June 21–25, 2004, Providence, RI: American Mathematical Society (AMS), 2007, pp. 101–123.

- [6] J. BERESTYCKI, E. BRUNET, S. C. HARRIS, AND M. ROBERTS, Vanishing corrections for the position in a linear model of FKPP fronts, Comm. Math. Phys., 349 (2017), pp. 857–893.
- [7] C. BESSE, G. FAYE, J.-M. ROQUEJOFFRE, AND M. ZHANG, The logarithmic Bramson correction for Fisher-KPP equations on the lattice Z, Trans. Am. Math. Soc., 376 (2023), pp. 8553–8619.
- [8] E. BOUIN AND C. HENDERSON, The Bramson delay in a Fisher-KPP equation with log-singular nonlinearity, Nonlinear Anal., 213 (2021), pp. Paper No. 112508, 30.
- [9] E. BOUIN, C. HENDERSON, AND L. RYZHIK, The Bramson delay in the non-local Fisher-KPP equation, Ann. Inst. H. Poincaré Anal. Non Linéaire, 37 (2020), pp. 51–77.
- [10] N. BOUTILLON, Large deviations and the emergence of a logarithmic delay in a nonlocal linearised Fisher-KPP equation, Nonlinear Analysis, 240 (2024), p. 113465.
- M. BRAMSON, Convergence of solutions of the Kolmogorov equation to travelling waves, vol. 285, American Mathematical Soc., 1983.
- [12] M. D. BRAMSON, Maximal displacement of branching brownian motion, Communications on Pure and Applied Mathematics, 31 (1978), pp. 531–581.
- [13] E. BRUNET AND B. DERRIDA, Shift in the velocity of a front due to a cutoff, Physical Review E, 56 (1997), p. 2597.
- [14] U. EBERT AND W. VAN SAARLOOS, Front propagation into unstable states: universal algebraic convergence towards uniformly translating pulled fronts, Phys. D, 146 (2000), pp. 1–99.
- [15] R. A. FISHER, The wave of advance of advantageous genes, Ann. of Eugenics, 7 (1937), pp. 355–369.
- [16] T. GILETTI, Monostable pulled fronts and logarithmic drifts, NoDEA, Nonlinear Differ. Equ. Appl., 29 (2022), p. 42. Id/No 35.
- [17] C. GRAHAM, The Bramson correction for integro-differential Fisher-KPP equations, Commun. Math. Sci., 20 (2022), pp. 563–596.
- [18] F. HAMEL AND G. NADIN, Spreading properties and complex dynamics for monostable reaction-diffusion equations, Comm. Partial Differential Equations, 37 (2012), pp. 511– 537.
- [19] F. HAMEL, J. NOLEN, J.-M. ROQUEJOFFRE, AND L. RYZHIK, A short proof of the logarithmic Bramson correction in Fisher-KPP equations, Netw. Heterog. Media, 8 (2013), pp. 275–289.
- [20] F. HAMEL AND L. ROQUES, Fast propagation for KPP equations with slowly decaying initial conditions, J. Differential Equations, 249 (2010), pp. 1726–1745.
- [21] A. HARAUX AND F. B. WEISSLER, Nonuniqueness for a semilinear initial value problem, Indiana Univ. Math. J., 31 (1982), pp. 167–189.

- [22] A. N. KOLMOGOROV, I. G. PETROVSKY, AND N. S. PISKUNOV, Etude de l'équation de la diffusion avec croissance de la quantité de matière et son application à un problème biologique, Bull. Univ. Etat Moscou, Sér. Inter. A 1 (1937), pp. 1–26.
- [23] K.-S. LAU, On the nonlinear diffusion equation of Kolmogorov, Petrovsky, and Piscounov, Journal of Differential Equations, 59 (1985), pp. 44–70.
- [24] J. NOLEN, J.-M. ROQUEJOFFRE, AND L. RYZHIK, Convergence to a single wave in the Fisher-KPP equation, Chin. Ann. Math., Ser. B, 38 (2017), pp. 629–646.
- [25] —, Refined long-time asymptotics for Fisher-KPP fronts, Commun. Contemp. Math., 21 (2019), p. 25. Id/No 1850072.
- [26] F. ROTHE, Convergence to travelling fronts in semilinear parabolic equations, Proc. R. Soc. Edinb., Sect. A, Math., 80 (1978), pp. 213–234.
- [27] K. UCHIYAMA, The behavior of solutions of some nonlinear diffusion equations for large time, J. Math. Kyoto Univ., 18 (1978), pp. 453–508.
- [28] W. VAN SAARLOOS, Front propagation into unstable states, Physics reports, 386 (2003), pp. 29–222.